

Transport through a nano-wire in a magnetic field

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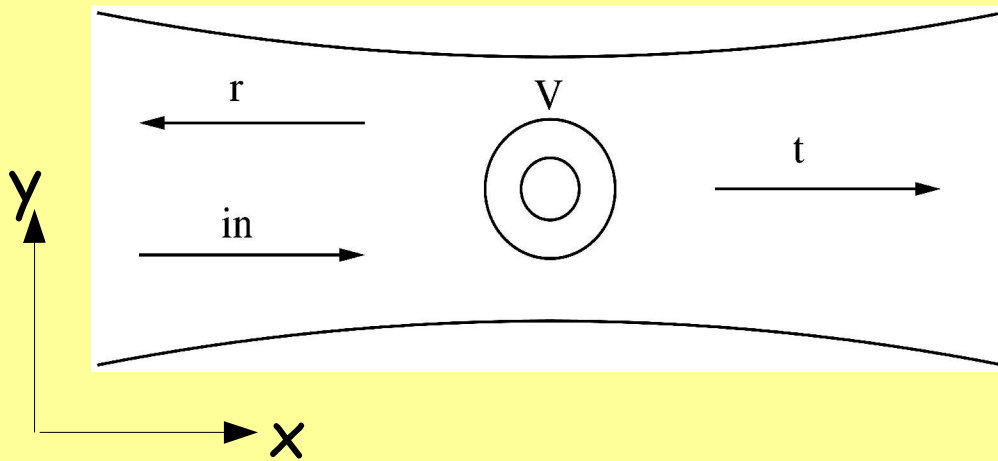
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Guðný Guðmundsdóttir, Ingibjörg Magnúsdóttir, Jens Hjörleifur Bárðarson, Viðar Guðmundsson, Andrei Manolescu, Chi-Shung Tang, and Yu-Yu Lin

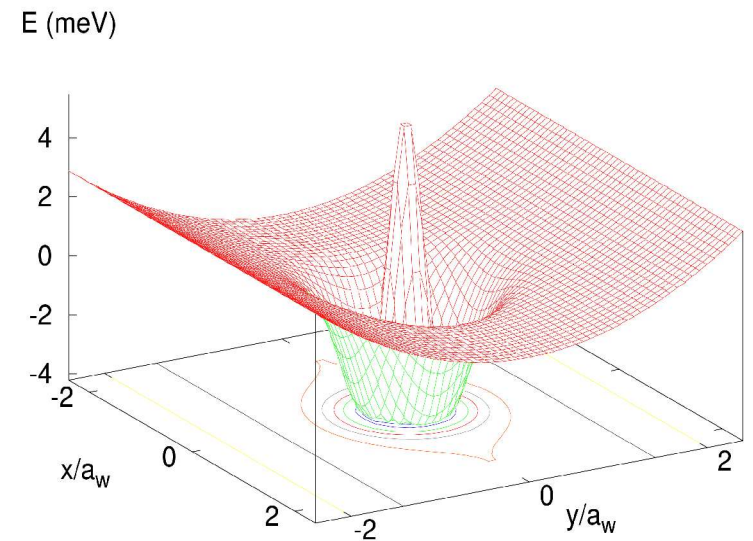
Quantum ring or dot embedded in wire

$B=B_z$, magnetic field

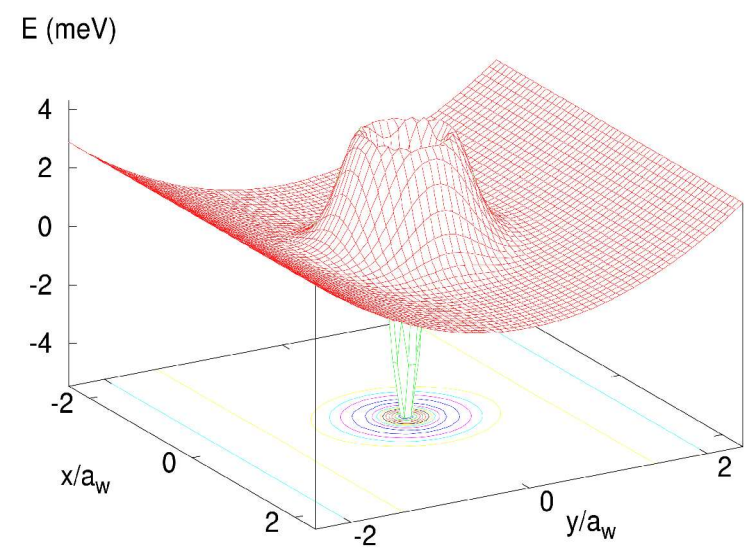


$$B=0.2 \text{ T} \rightarrow a_w = 32.8 \text{ nm}$$

$$E_0 = 1.0 \text{ meV}, V_1 = -12 \text{ meV}, V_2 = 18 \text{ meV}$$



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Model

$$H_0 = \frac{\hbar^2}{2m^*} \left[-i\nabla - \frac{eB}{\hbar c} y \hat{x} \right]^2 + V_c(y)$$

Center coordinate

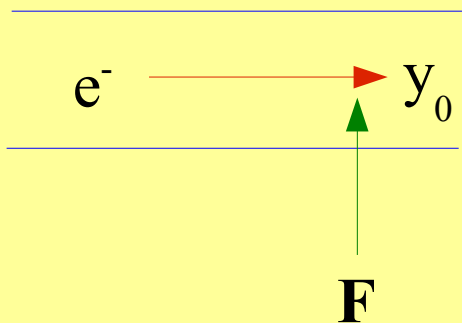
$$\psi^+(x, y, k_n) = e^{ik_n x} \chi_n(y - y_0)$$

$$E = \left(n + \frac{1}{2} \right) \hbar \Omega_w + \mathcal{K}(k_n)$$

Asymptotic region

$$\Omega_w = \sqrt{\omega_c^2 + \Omega_0^2}, \quad y_0 = k_n a_w^2 \frac{\omega_c}{\Omega_w}, \quad \omega_c = \frac{eB}{m^* c}$$

$$\mathcal{K}(k_n) = \frac{(k_n a_w)^2}{2} \left(\frac{\hbar^2 \Omega_0^2}{\hbar \Omega_w} \right), \quad a_w^2 \Omega_w = \frac{\hbar}{m^*}$$



B → "Non-orthogonality"

Solution

Mixed momentum-
coordinate representation

$$\Psi_E(p, y) = \int dx \psi_E(x, y) e^{-ipx}$$

$$\Psi_E(p, y) = \sum_n \varphi_n(p) \phi_n(p, y)$$

Expansion in terms of eigenfunctions of the shifted harmonic oscillator -> transport mode "n"

...transforms the Schrödinger eq.

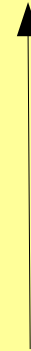
$$\kappa_n(q)\varphi_n(q) + \sum_{n'} \int \frac{dp}{2\pi} V_{nn'}(q, p)\varphi_{n'}(p) = (E - E_n)\varphi_n(q)$$

$$V_{nn'}(q, p) = \int dy \phi_n^*(q, y)V(q - p, y)\phi_{n'}(p, y)$$

$$V(q - p, y) = \int dx e^{-i(q-p)x}V_{sc}(x, y)$$

a set of coupled integral eq.'s

Scattering potential



...rewrite

Non-local

$$\left[-(qa_w)^2 + (k_n(E)a_w)^2 \right] \varphi_n(q) = \frac{2\hbar\Omega_w}{(\hbar\Omega_0)^2} \sum_{n'} \int \frac{dp}{2\pi} V_{nn'}(q, p) \varphi_{n'}(p)$$

Effective band
momentum

$$(E - E_n) = \frac{[k_n(E)]^2 (\hbar\Omega_0)^2}{2 \hbar\Omega_w}$$

$$\left[-(qa_w)^2 + (k_n(E)a_w)^2 \right] \varphi_n^0(q) = 0$$

Free equation

Suggests an interpretation....

...a Green's function

$$\left[-(qa_w)^2 + (k_n(E)a_w)^2 \right] G_E^n(q) = 1$$

Lippmann-
Schwinger
eq. in q-space

$$\varphi_n(q) = \varphi_n^0(q) + G_E^n(q) \sum_{n'} \int \frac{dp a_w}{2\pi} \tilde{V}_{nn'}(q, p) \varphi_{n'}(p)$$

$$\varphi = \varphi^0 + G\tilde{V}\varphi^0 + G\tilde{V}G\tilde{V}\varphi^0 + \dots = (1 + G\tilde{T})\varphi^0$$

$$\tilde{T}_{nn'}(q, p) = \tilde{V}_{nn'}(q, p) + \sum_{m'} \int \frac{dk a_w}{2\pi} \tilde{V}_{nm'}(q, k) G_E^{m'}(k) \tilde{T}_{nm'}(k, p).$$

transformed into eq.'s for T-matrix

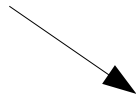
Supplies

Conductance



$$G(E) = \frac{2e^2}{h} \text{Tr}[\mathbf{t}^\dagger(E)\mathbf{t}(E)]$$

Transmission
amplitudes



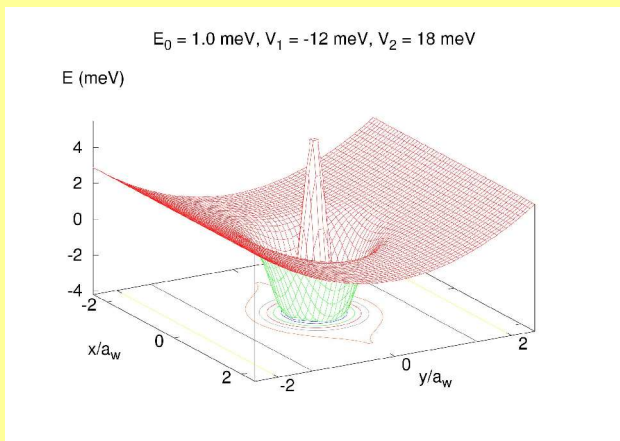
$$t_{nm}(E) = \delta_{nm} - \frac{i\sqrt{(k_m/k_n)}}{2(k_m a_w)} \left(\frac{\hbar\Omega_0}{\hbar\Omega_w} \right)^2 \tilde{T}_{nm}(k_n, k_m)$$

$$\psi_E(x, y) = e^{ik_n x} \phi_n(k_n, y) + \sum_m \int \frac{dq a_w}{2\pi} e^{iqx} G_E^m(q) \tilde{T}_{mn}(q, k_n) \phi_m(q, y)$$

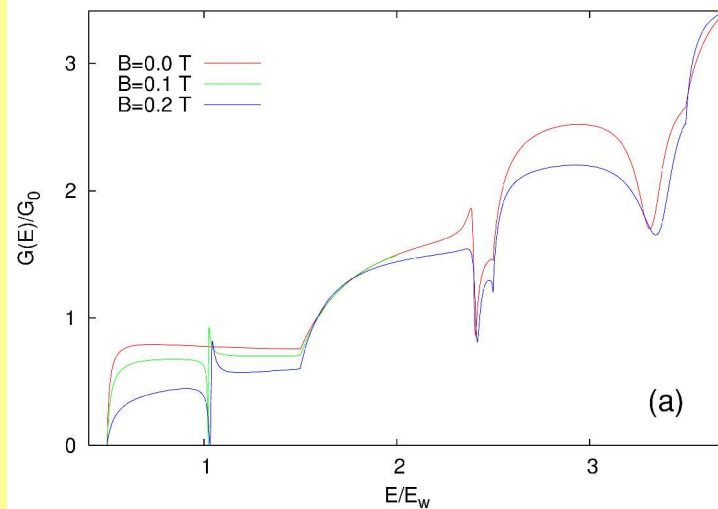
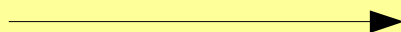


Wave functions

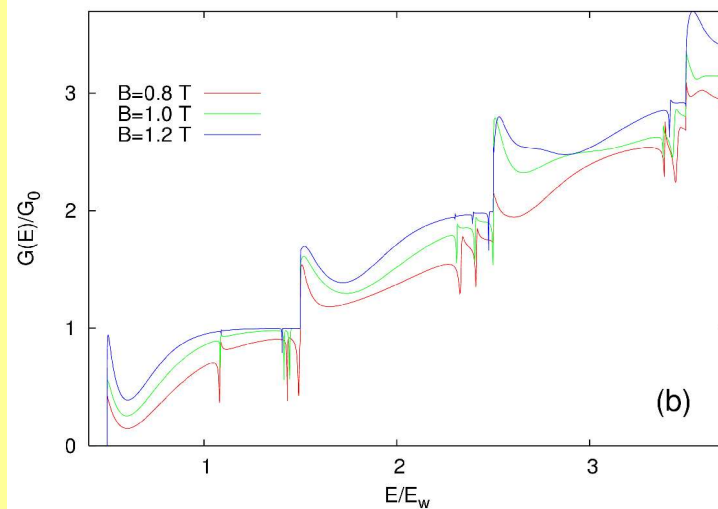
Quantum ring



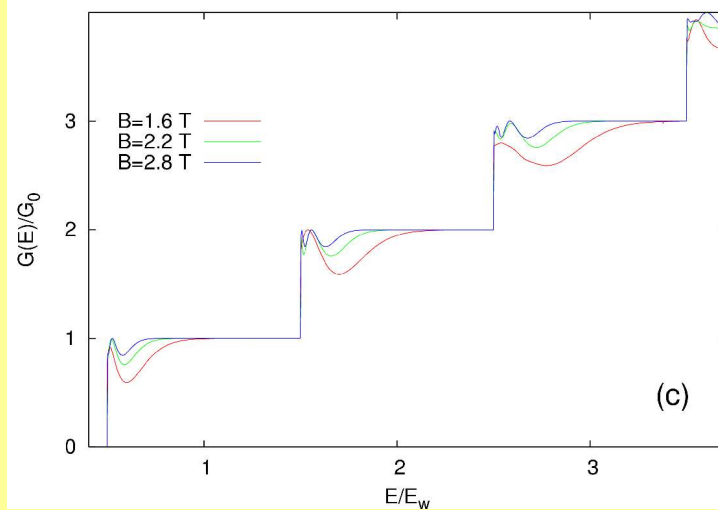
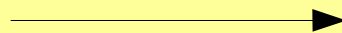
Low B



Intermediate B

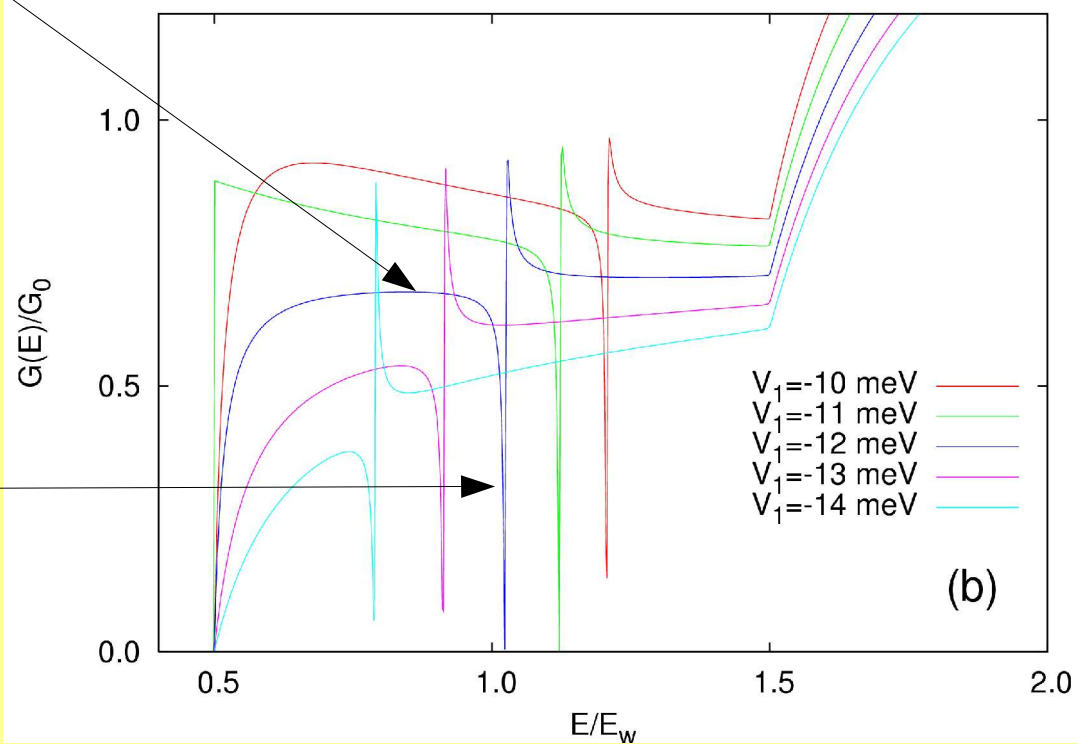
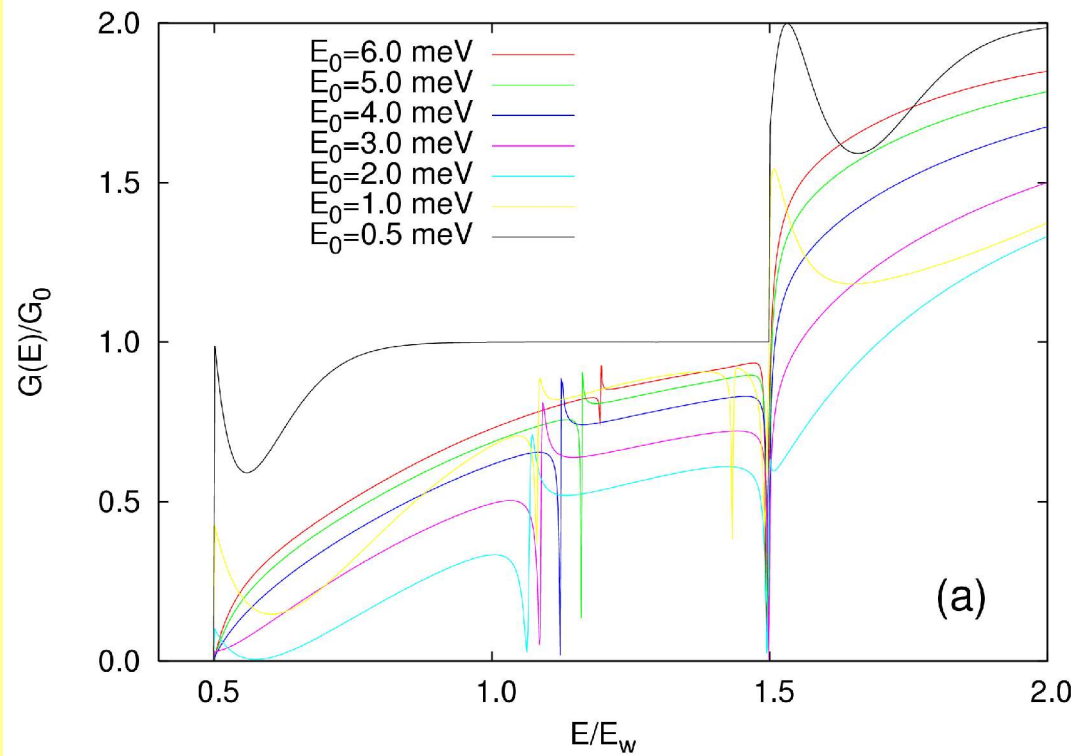
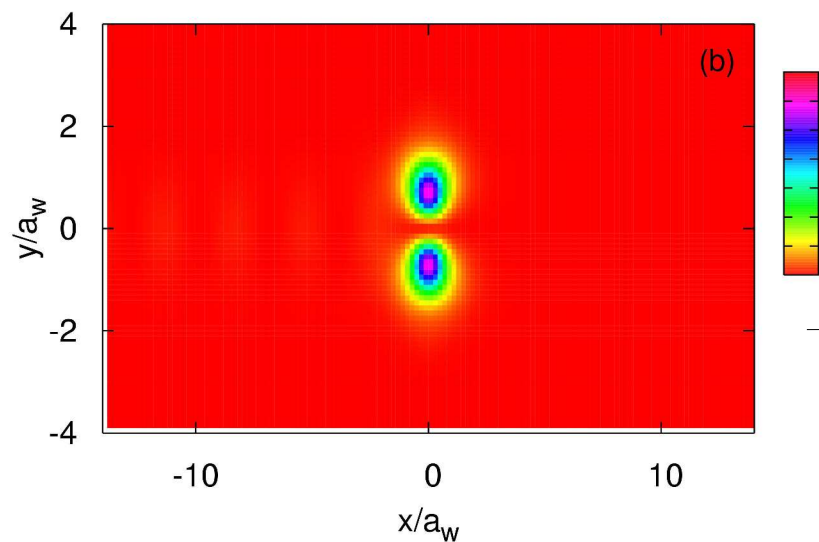
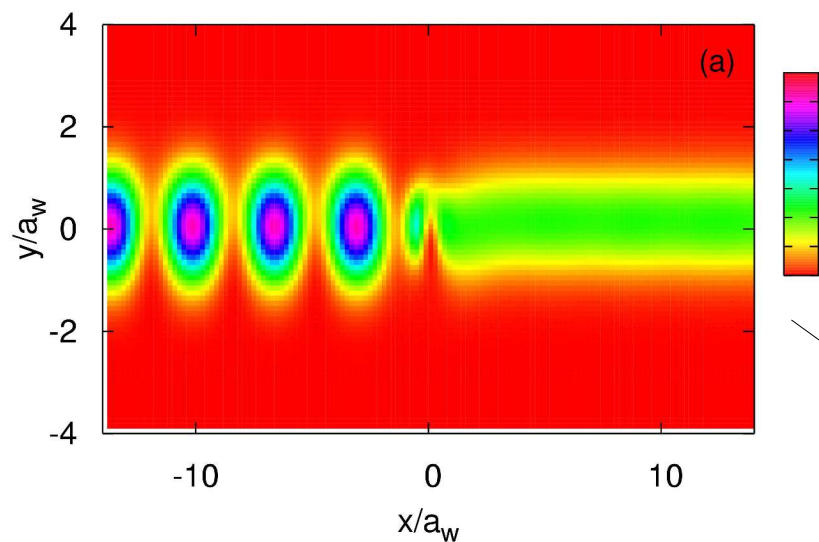


High B

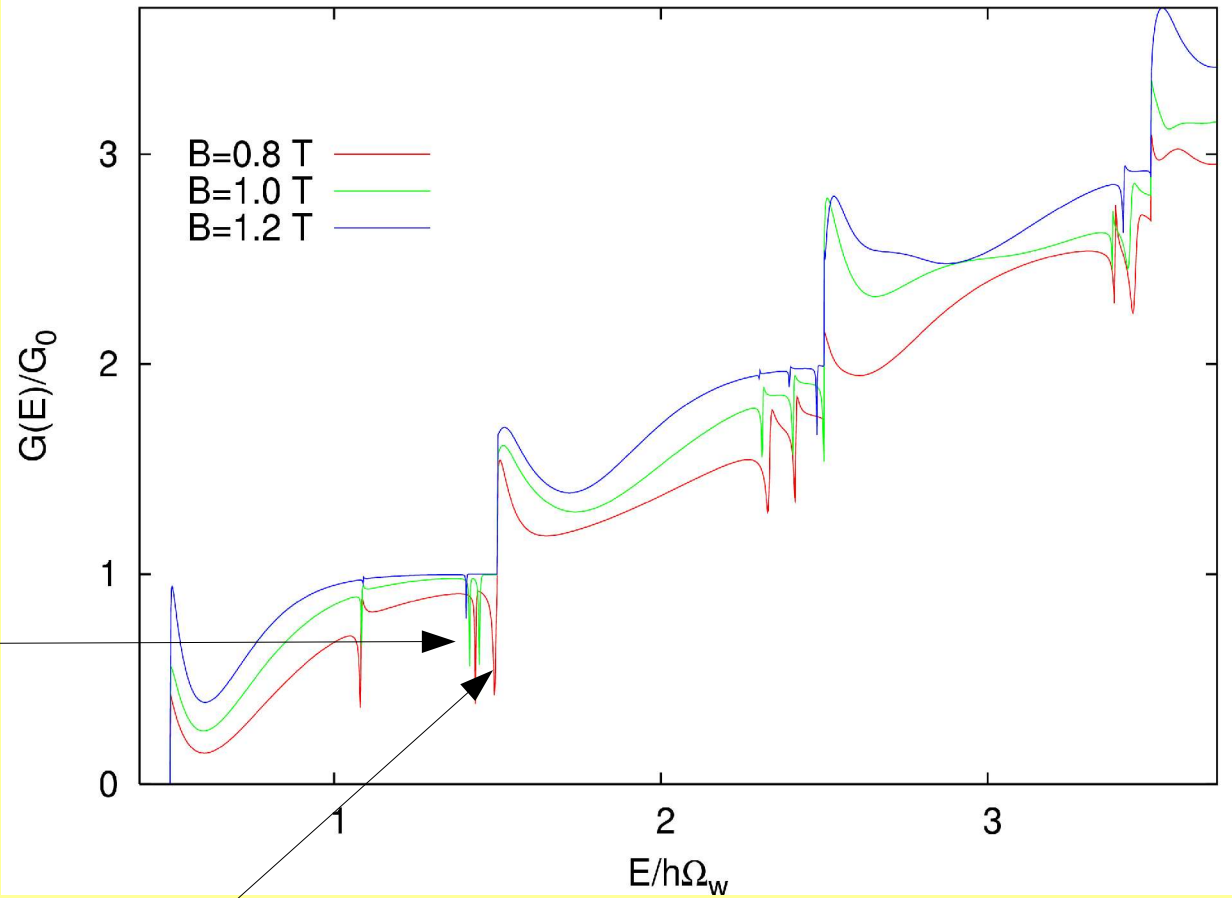
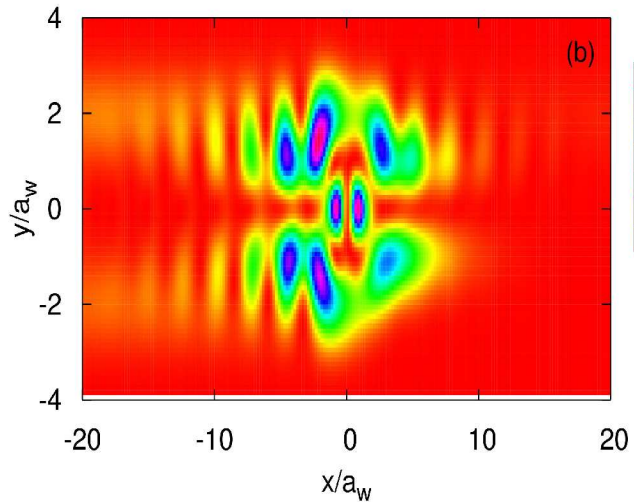
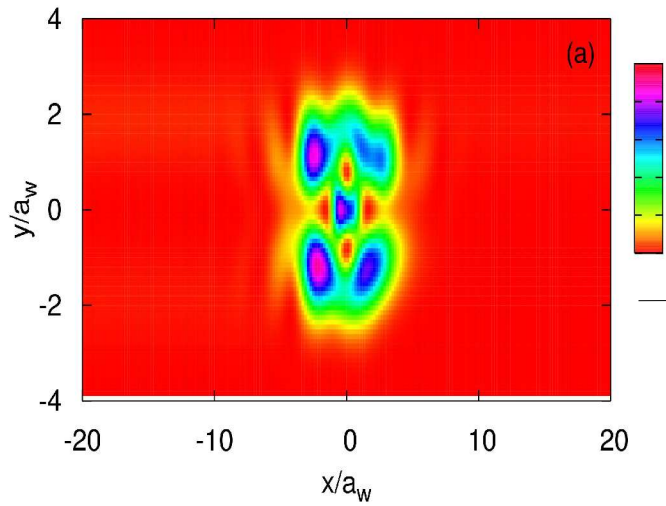


Fano resonance

(Forbidden at $B=0$)



Evanescent states

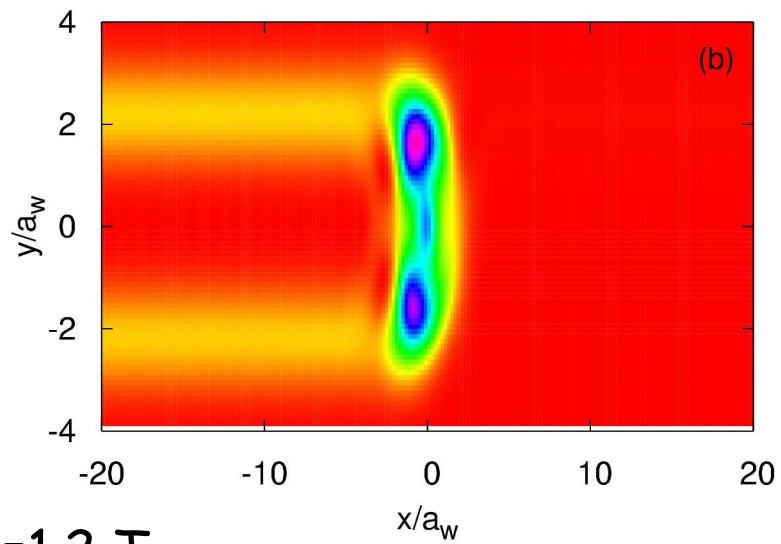
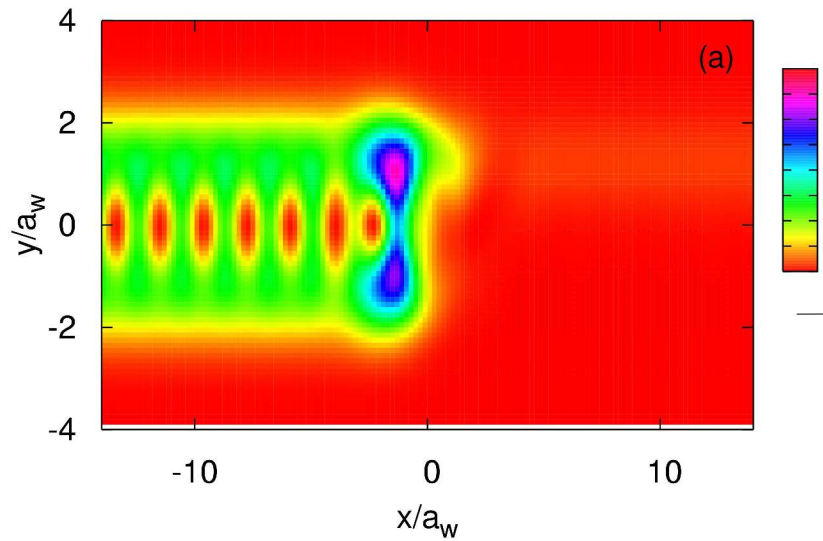


Partial separation of
in and reflected waves

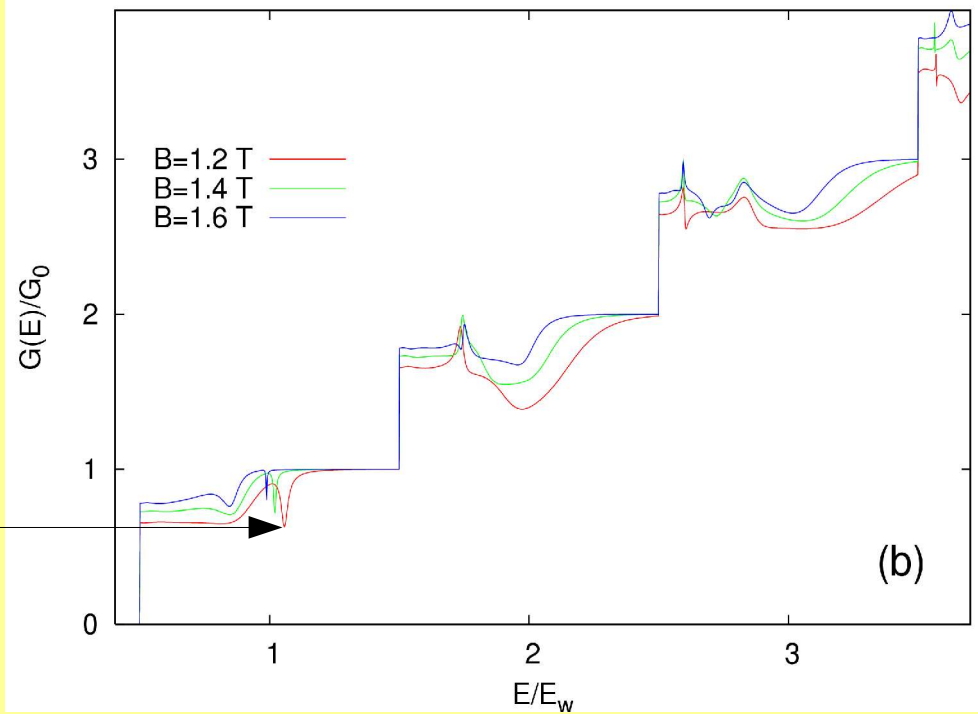
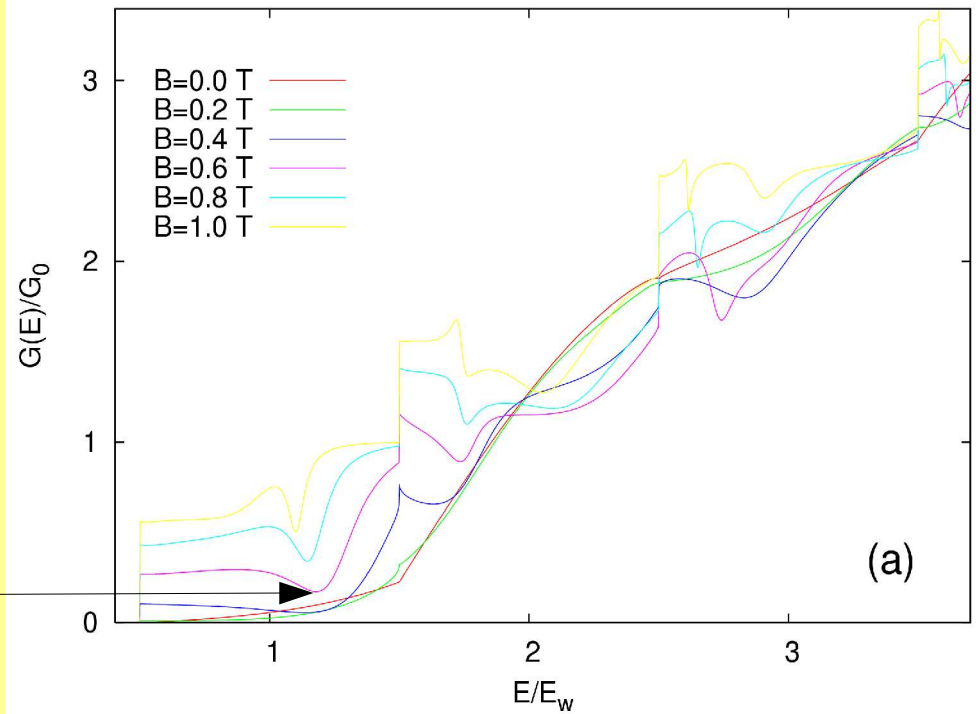
$B=0.8$ T

Embedded dot

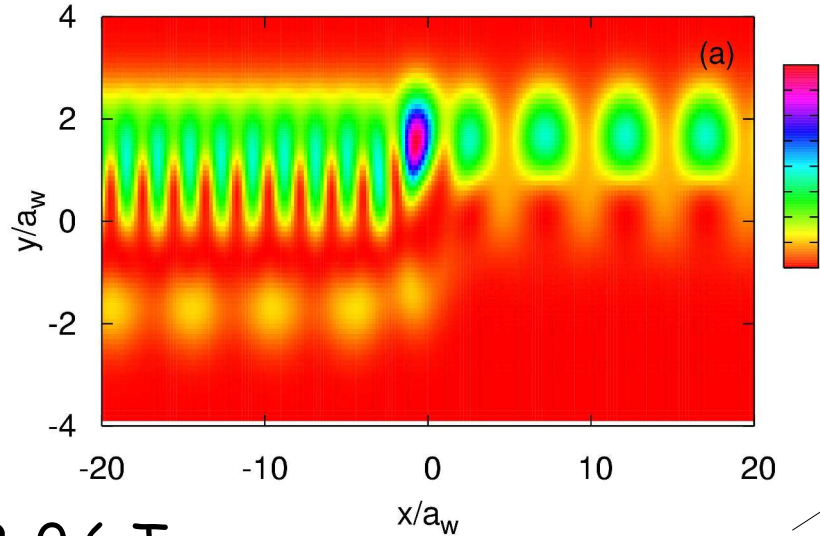
$B=0.6$ T



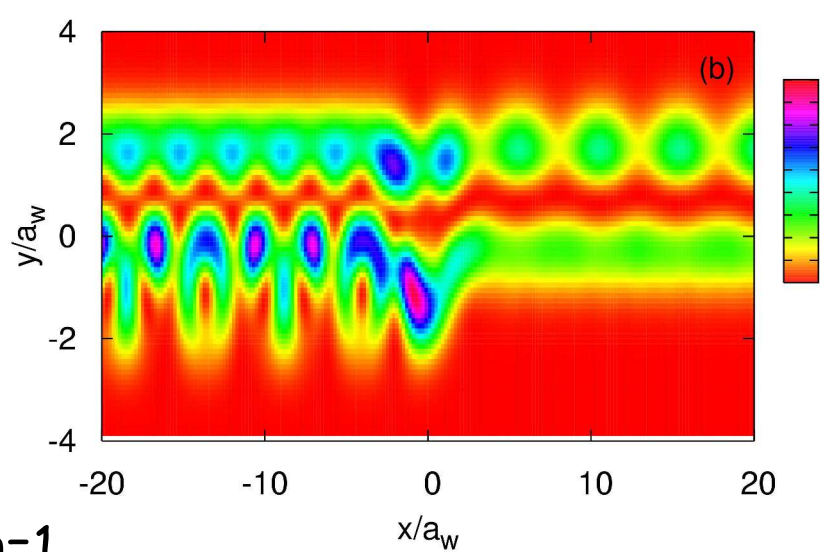
$B=1.2$ T



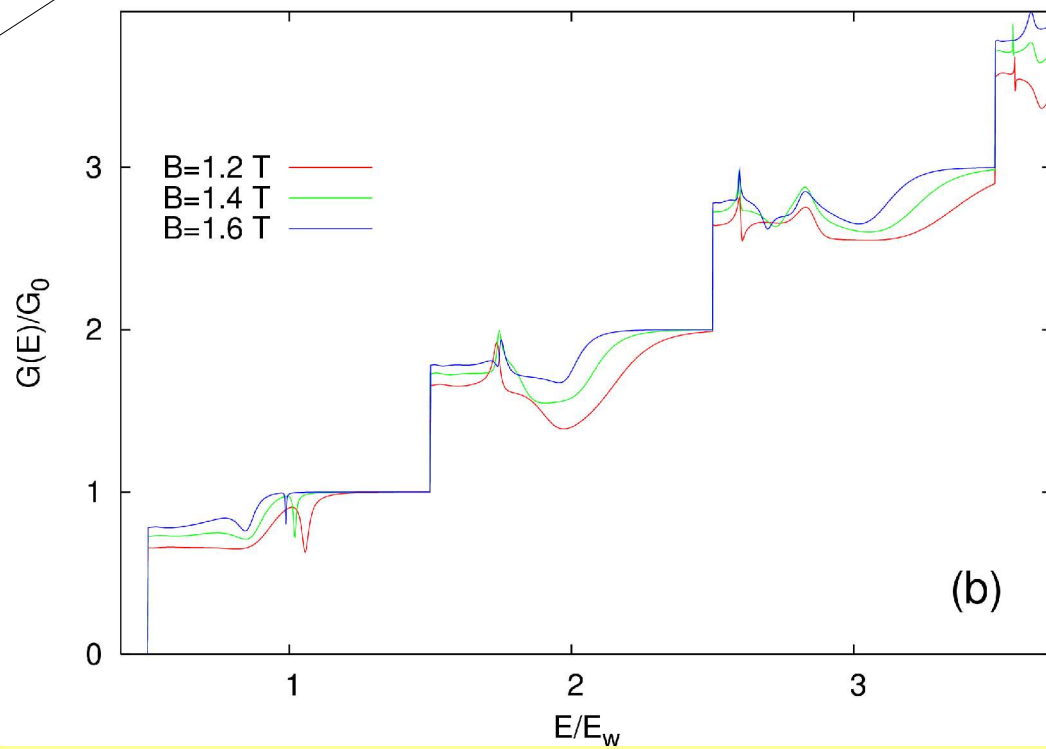
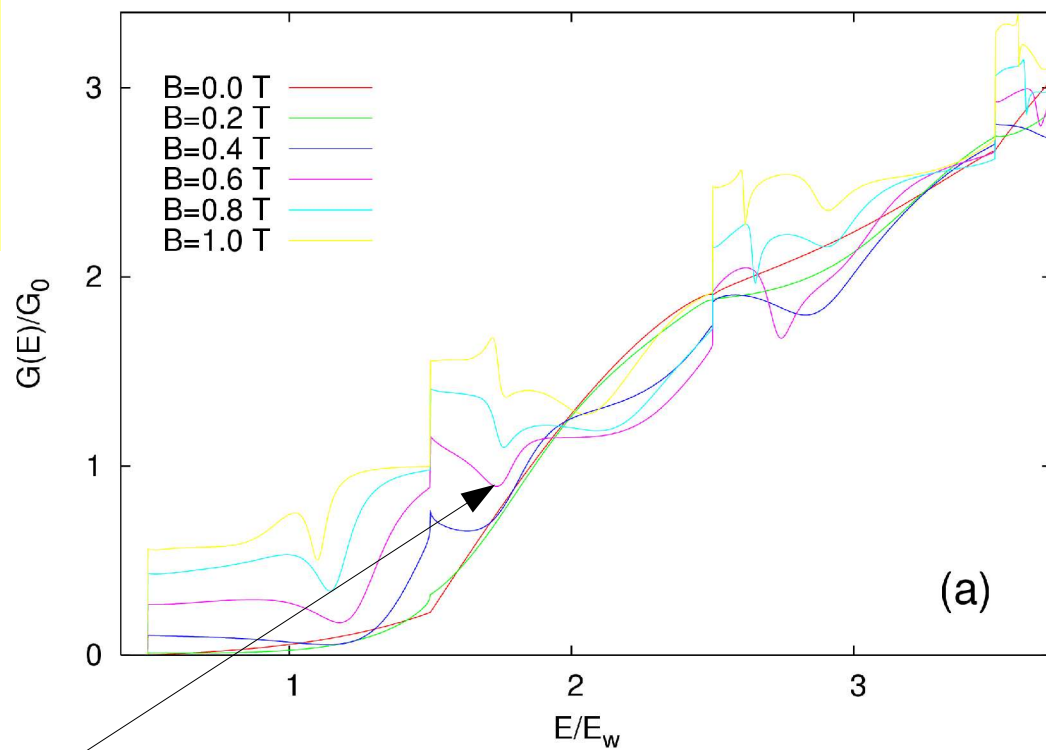
$n=0$



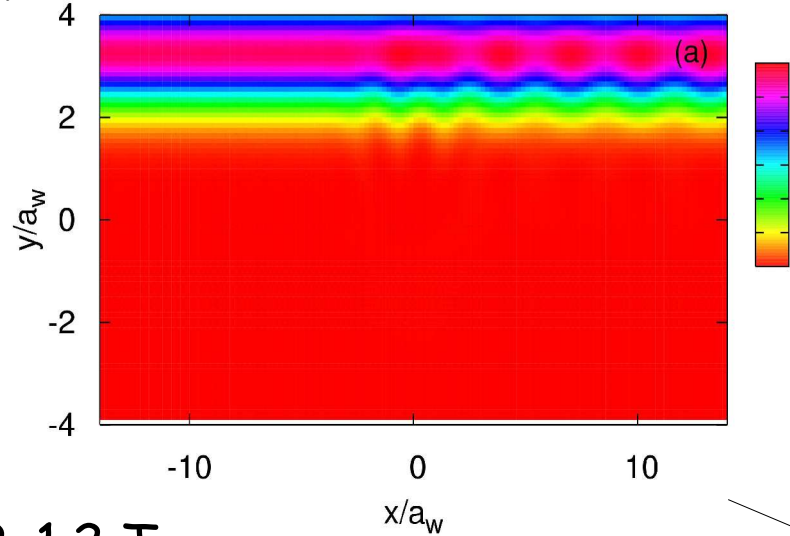
$B=0.6$ T



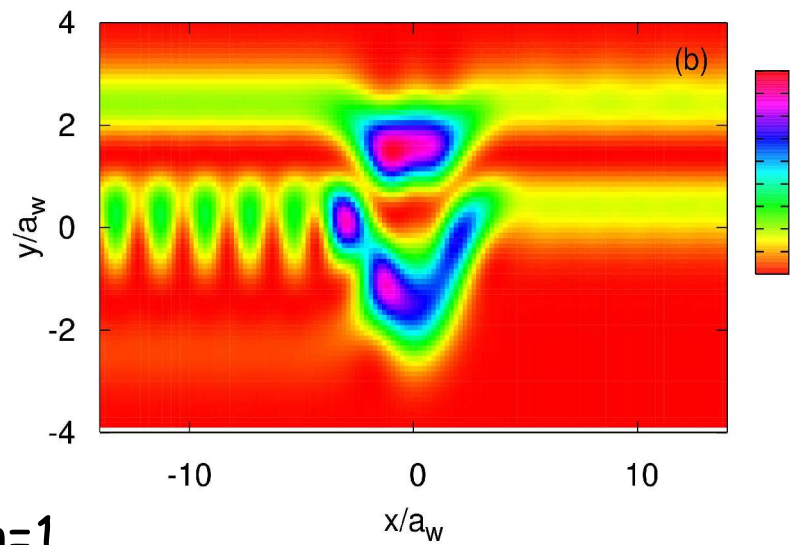
$n=1$



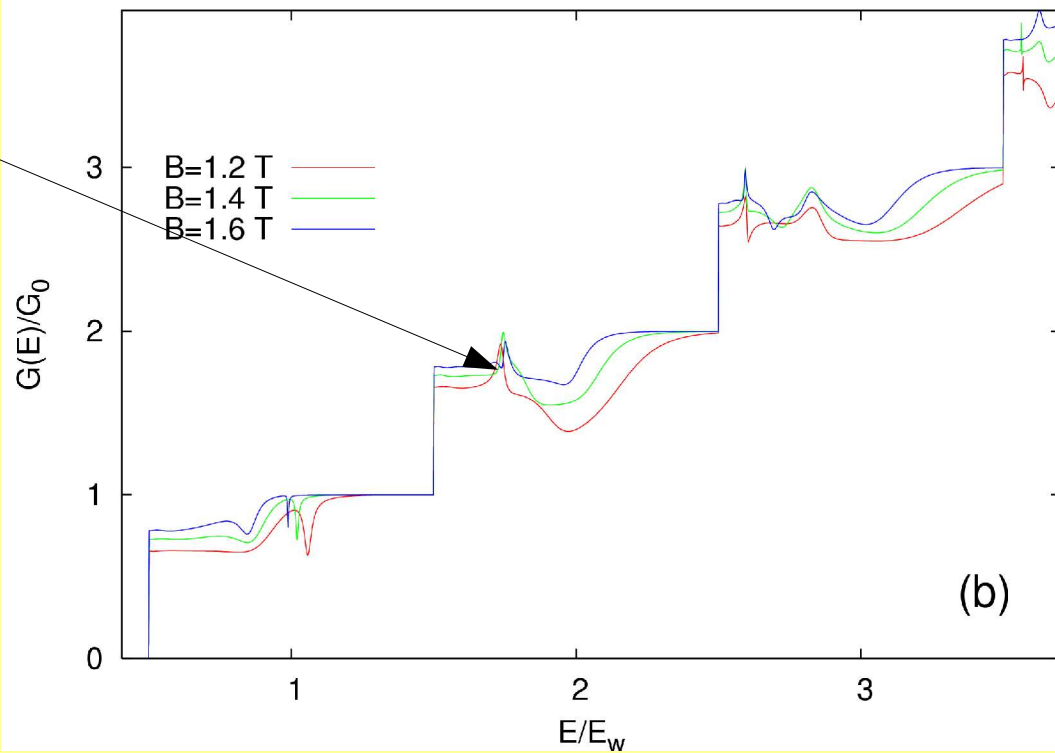
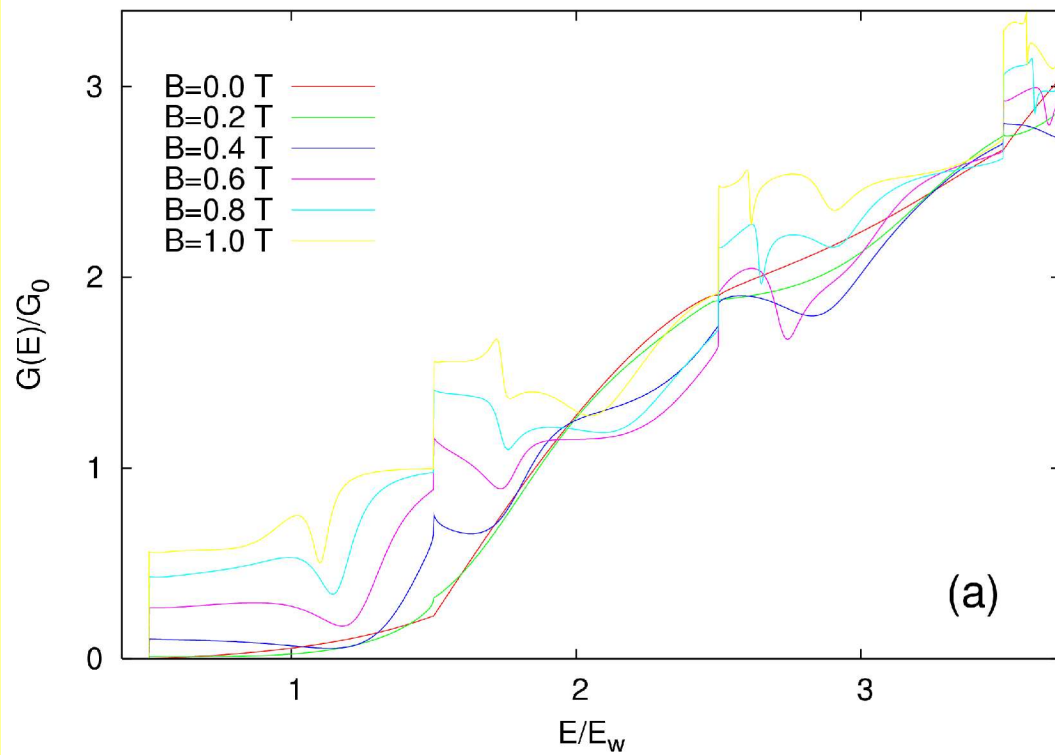
$n=0$



$B=1.2$ T

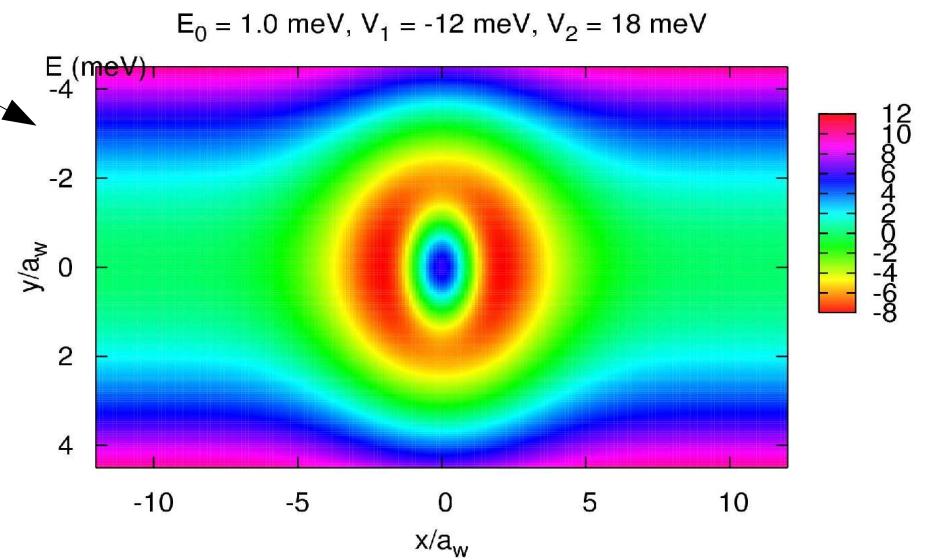
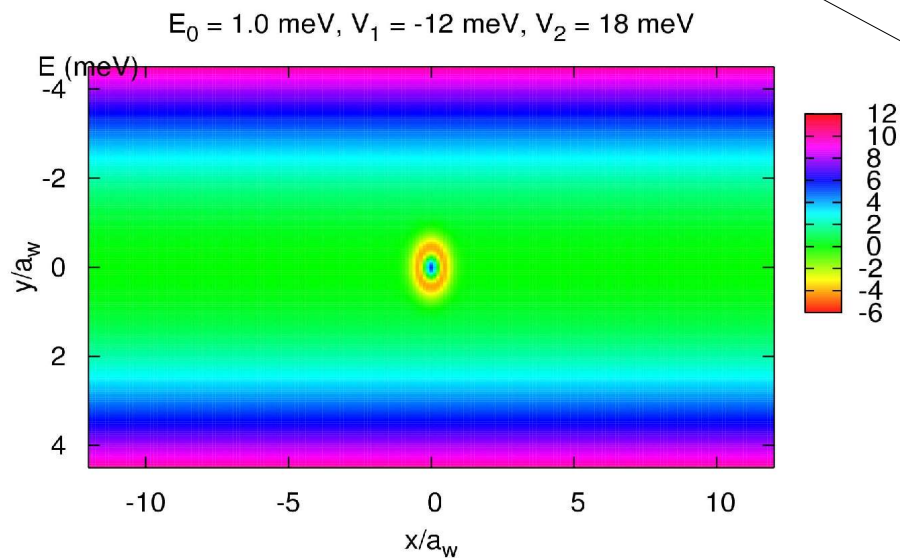
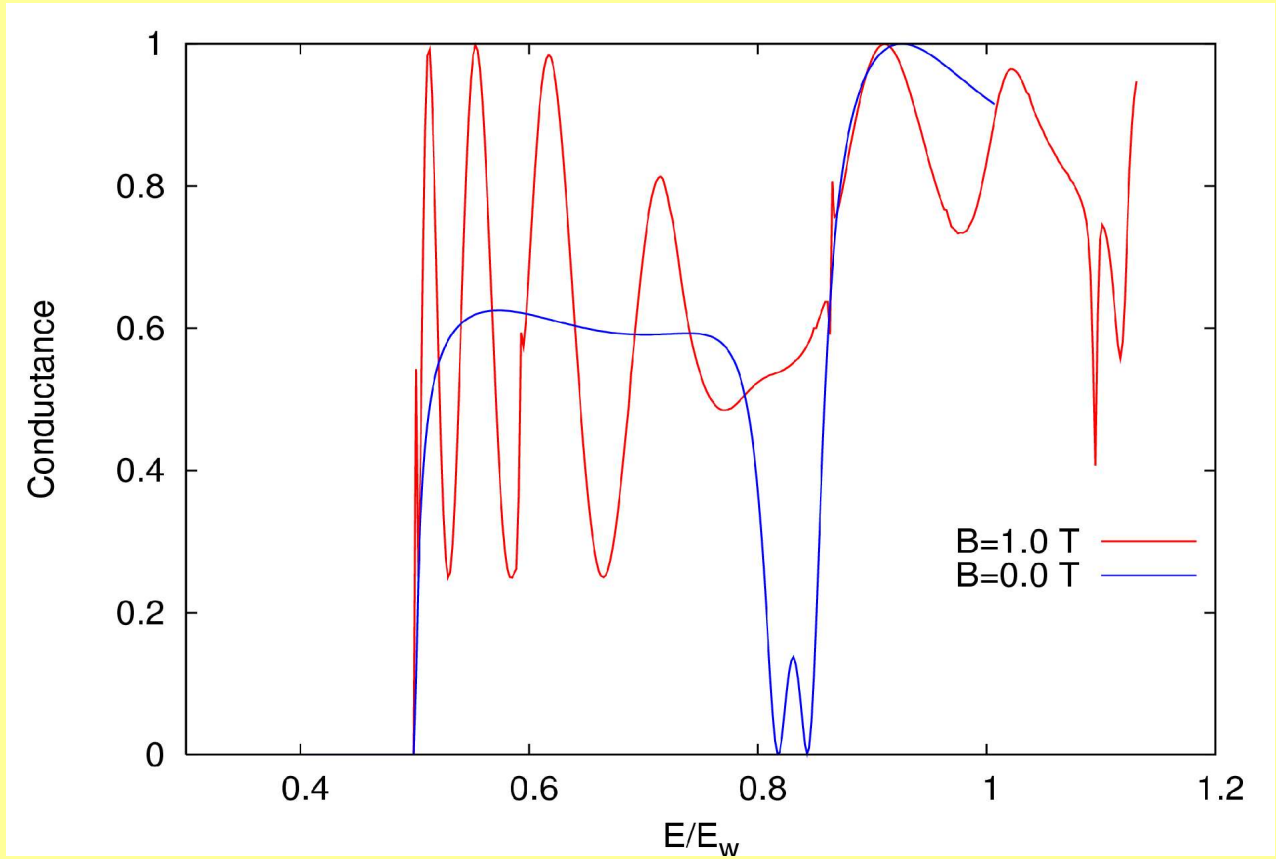


$n=1$



Aharonov-Bohm

What happens in a **larger** ring



We can describe scattering in magnetic field

Relaxation of selection rules

Edge-modes, channels

Aharonov-Bohm oscillations

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