

# *Time-dependent magnetotransport in a quantum wire*

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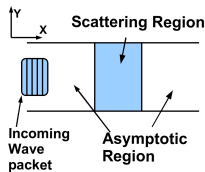
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Tainan, September, 2007

# Two cases of time-dependent magnetotransport

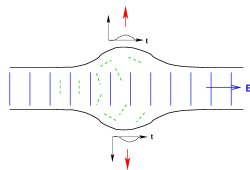
## Wave packet transport

- Static potential
- Elastic scattering
- Life-time of quasi-bound states and resonances
- Delay times



## Modulation of a current

- Plane in-wave  $\rightarrow$  sharp in-energy  $E$
- Time-dependent potential
- Potential flashed smoothly on and off, not periodic
- Inelastic scattering



## Asymptotic regions, both cases

Landau gauge:  $\mathbf{A} = -By\hat{\mathbf{x}} \rightarrow \mathbf{B} = B\hat{\mathbf{z}}, \quad l = \sqrt{\hbar c/(eB)}$

Equation of motion

$$i\hbar\partial_t\Psi(\mathbf{r}, t) = \left\{ -\frac{\hbar^2}{2m^*} \left( \nabla^2 - \frac{2i}{l^2}y\partial_x - \frac{y^2}{l^4} \right) + \frac{1}{2}m^*\Omega_0^2y^2 \right\} \Psi(\mathbf{r}, t)$$

Fourier transform:  $(x, y, t) \rightarrow (p, y, \omega)$

$$\Psi(\mathbf{r}, t) = \int \frac{dp}{2\pi} \frac{d\omega'}{2\pi} e^{i(px - \omega't)} \Psi(p, y, \omega')$$

Separation into modes in  $(p, y)$ -space, (S. A. Gurvitz, PRB **51**, 7123 (1995))

$$\Psi(q, y, t) = \sum_n \varphi_n(q, t) \phi_n(q, y)$$

H.O. eigenfunctions  $\phi_n(q, y)$ , (not orthogonal for different  $q$ 's)

Frequency  $\Omega_w = \sqrt{\omega_c^2 + \Omega_0^2}$ , cyclotron frequency  $\hbar\omega_c = eB/(m^*c)$

Shifted by  $y_0 = qa_w^2\omega_c/\Omega_w$ , new length scale  $a_w = \sqrt{\hbar/(m^*\Omega_w)}$

For a **static** system:

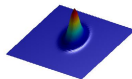
Energy bands

$$E_{nq} = \hbar\omega_{nq}^0 = E_n^0 + U_w \frac{(qa_w)^2}{2}, \quad U_w = \frac{(\hbar\Omega_0)^2}{\hbar\Omega_w}$$

With band edges  $E_n^0 = \hbar\Omega_w(n + 1/2)$  for  $n = 0, 1, 2, \dots$

# Wave packet propagation

Wave packet



$$\varphi_n^0(p, \omega) = 2\pi g_n(p) \delta \left[ \omega - \frac{E_{np}}{\hbar} \right]$$

$$g_n(p) = \delta_{nn'} \exp \left[ -\gamma(p - p_0)^2 \right]$$

and the Green function

$$G_n(p, \omega) = \frac{1}{(k_n(\omega) a_w)^2 - (p a_w)^2}$$

with

$$k_n(\omega) a_w = \sqrt{\frac{(\hbar\omega - E_n^0)}{U_w}}$$

gives a Lippmann-Schwinger equation

$$\varphi_n(p, \omega) = \varphi_n^0(p, \omega) + G_n(p, \omega) \int_{-\infty}^{\infty} \frac{dq a_w}{2\pi} U_{nn'}(p, q) \varphi_{n'}(q),$$

where

$$U_{nn'}(p, q) U_w = \int_{-\infty}^{\infty} dy \phi_{n'}^*(q, y) V_{sc}(p - q, y) \phi_n(p, y)$$

With a  $T$ -matrix

$$T_{nn'}(p, q, \omega) = U_{nn'}(p, q) + \sum_m \int_{-\infty}^{\infty} \frac{dk a_w}{2\pi} U_{nm}(p, k) G_m(k, \omega) T_{mn'}(k, q, \omega)$$

the wave function

$$\Psi(x, y, t) = \Psi_0(x, y, t) + \Psi_{\text{sc}}(x, y, t)$$

with the in-wave

$$\Psi_0(x, y, t) = \sum_n \int_{-\infty}^{\infty} dp g_n(p) \phi_n(p, y) e^{i(px - \omega_{np}^0 t)}$$

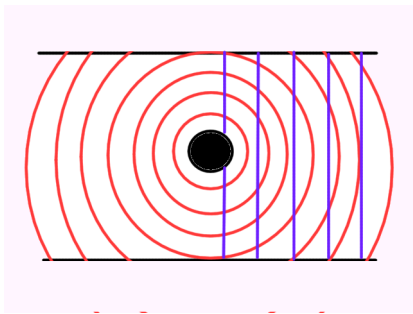
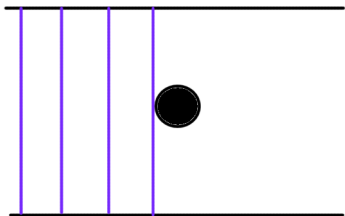
becomes

$$\begin{aligned} \Psi_{\text{sc}}(x, y, t) &= \sum_{n'} \int_{E_{n'}^0/\hbar}^{\infty} d\omega e^{-i\omega t} \frac{\Omega_w g_{n'}[k_{n'}(\omega)]}{\Omega_0^2 |k_n'(\omega) a_w|} \\ &\times \sum_n \int_{-\infty}^{\infty} \frac{dp a_w}{2\pi} G_n(p, \omega) e^{ipx} T_{nn'}(p, k_{n'}(\omega)) \phi_n(p, y) \end{aligned}$$

... or graphically

$$\Psi(x, y, t) = \Psi_0(x, y, t) + \Psi_{sc}(x, y, t).$$

before and after scattering

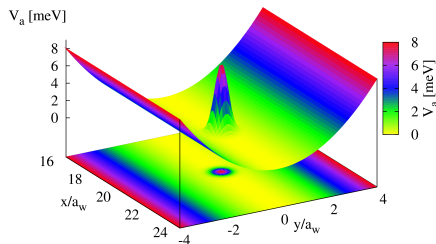




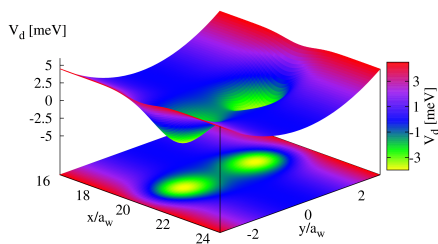
# Propagation of a wave packet

## Static potentials

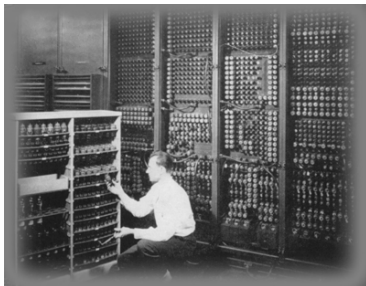
### Antidot



### Parallel double dot



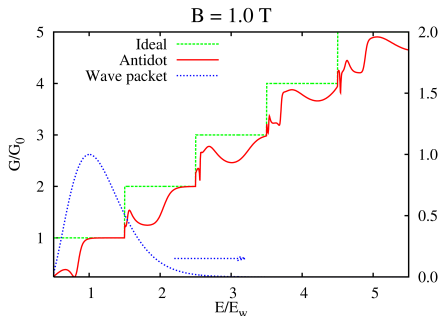
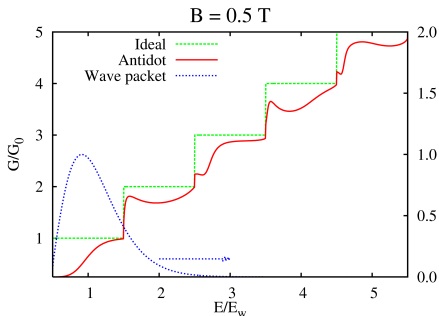
- Calculations implemented on the cluster jotunn.rhi.hi.is
- $\omega$ -integration programmed for parallel execution



- Large calculations:
  - 20 nodes  $\rightarrow$   $\sim$  40 hours
  - 3 GB of memory used on each node

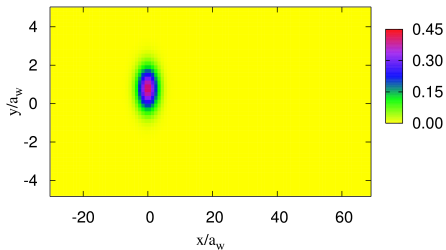
# Antidot

## Static conductance – wave packet

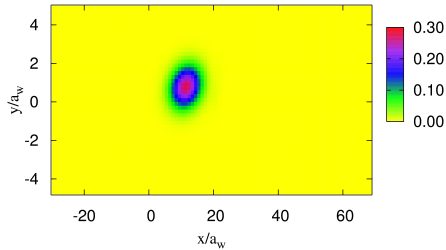


# Antidot, $B = 0.5 \text{ T}$

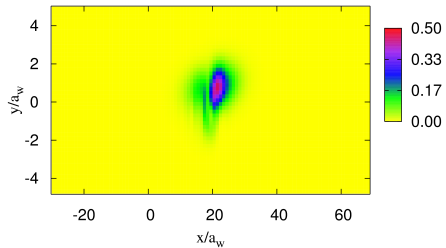
$t = 0 \text{ ps}$



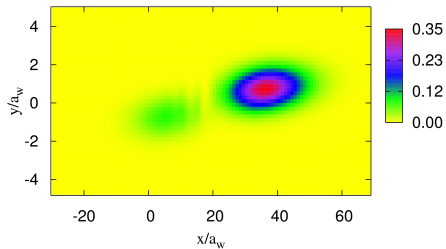
$t = 8 \text{ ps}$



$t = 15 \text{ ps}$

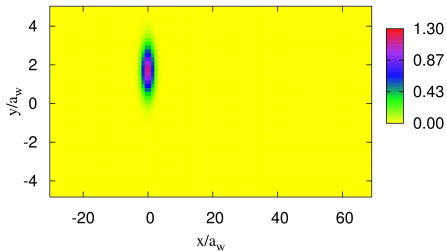


$t = 28 \text{ ps}$

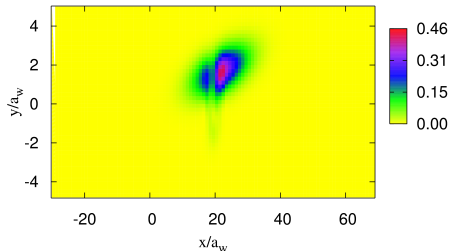


# Antidot, $B = 1.0$ T

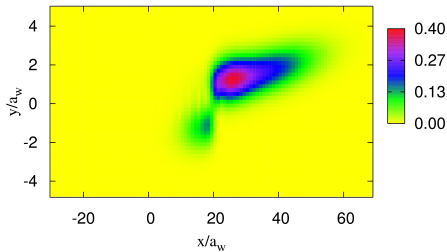
$t = 0$  ps



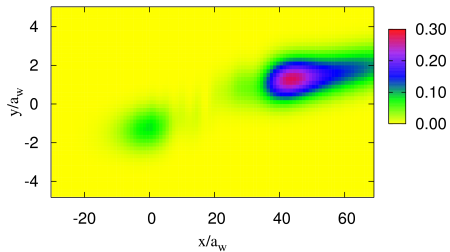
$t = 15$  ps



$t = 25$  ps

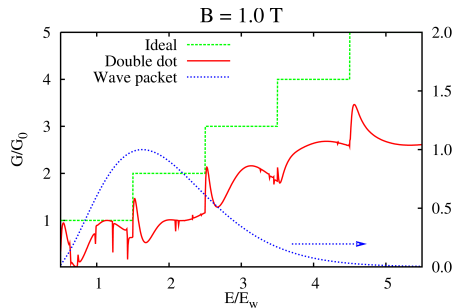
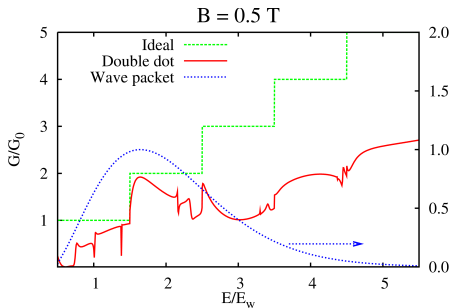


$t = 40$  ps



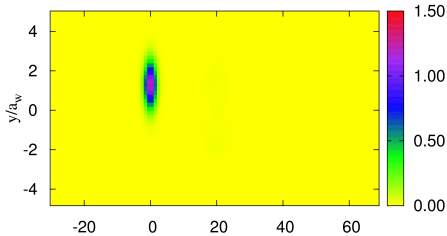
## Parallel double dot

### Static conductance – wave packet

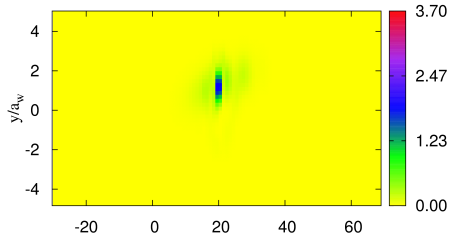


# Parallel double dot, $B = 0.5 T$

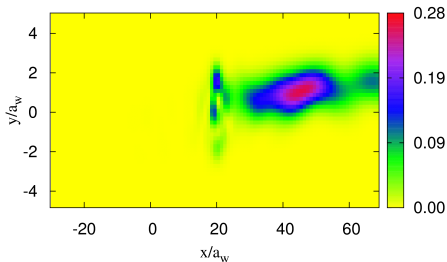
$t = 0$  ps



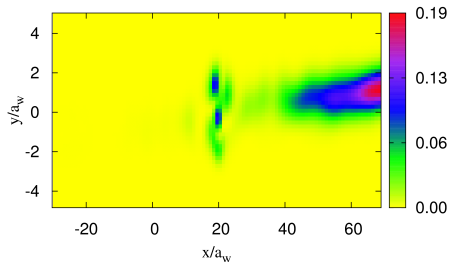
$t = 9$  ps



$t = 25$  ps

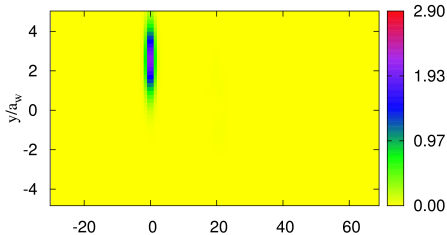


$t = 38$  ps

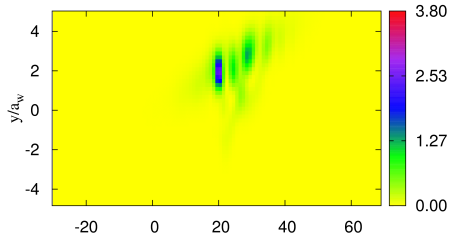


# Parallel double dot, $B = 1.0 \text{ T}$

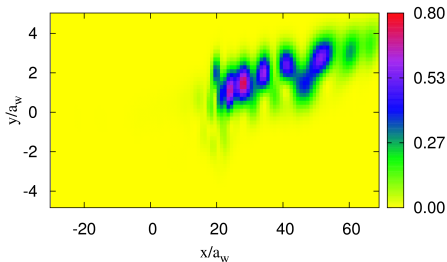
$t = 0 \text{ ps}$



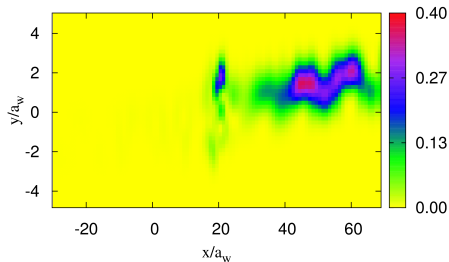
$t = 11 \text{ ps}$



$t = 21 \text{ ps}$



$t = 40 \text{ ps}$

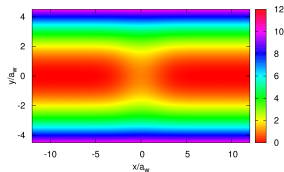
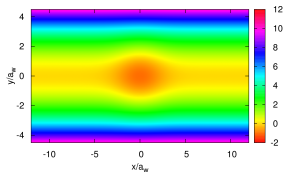
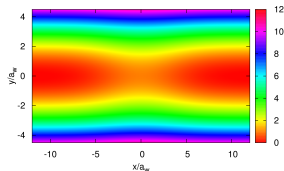
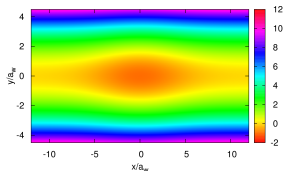




# Current modulation

$$V_{\text{sc}}(\mathbf{r}, t) = V_0 e^{-\beta r^2} e^{-\gamma t} \cos(\Omega t),$$

view at  $t = 0$ :



$V_0 = \pm 1.0 \text{ meV}$ ,  $\Omega = 0.2\Omega_w$ ,  $\gamma = 1.0\Omega_w^2$ ,  $\beta = 1 \text{ or } 4 \times 10^{-4} \text{ nm}^{-2}$ ,  $\rightarrow$  **one smooth flash**

Plane in-wave

$$\varphi_m^0(q, \omega) = (2\pi)^2 \delta(q - k_n) \delta(\omega - \omega_{nq}^0) \delta_{m,n}$$

Green function

$$\{\hbar\omega - \hbar\omega_{nq}^0\} G_0^n(q, \omega) = 1$$

$T$ -matrix

$$T_{nn'}(q\omega, p\nu) = V_{nn'}^{\text{sc}}(q\omega, p\nu) + \sum_{m'} \int \frac{dk}{2\pi} \frac{d\omega'}{2\pi} V_{nm'}^{\text{sc}}(q\omega, k\omega') G_0^{m'}(k\omega') T_{m'n'}(k\omega', p\nu)$$

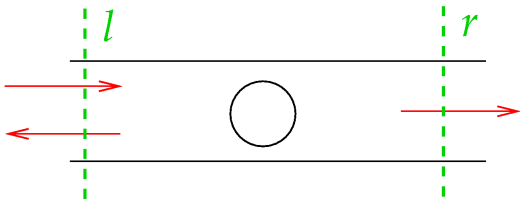
and full wave function

$$\Psi(\mathbf{r}, t) = e^{i(k_n x - \omega_{nk}^0 t)} \phi_n(k_n, y) + \sum_m \int \frac{dq}{2\pi} \frac{d\omega}{2\pi} e^{i(qx - \omega t)} G_0^m(q\omega) T_{mn}(q\omega, k_n \omega_{nk}^0) \phi_m(q, y)$$

Left and right current of state  $\alpha$

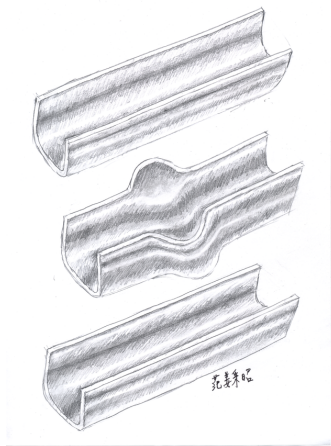
$$(I_{\alpha}^{r,l}(t))_x = \frac{\hbar}{m^*} \Re \left\{ \int_{-\infty}^{\infty} dy (\Psi_{\alpha}^{r,l})^* D_x \Psi_{\alpha}^{r,l} \right\}$$

with  $\hbar D_x = (p_x + (e/c)A_x) = \hbar(-i\partial_x - y/l^2)$



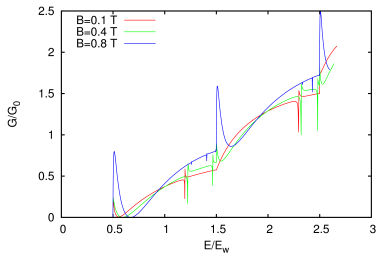
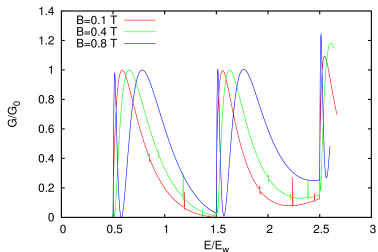
- Contributions from any point in sc-region for all earlier times
- Calculate for state  $\alpha$  at Fermi energy
- Inelastic, any outstate possible, evanescent states explicitly in  $G$

## Smooth well-like pulse

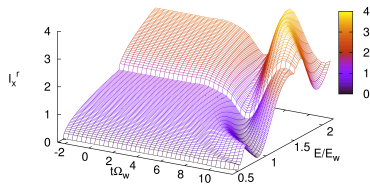
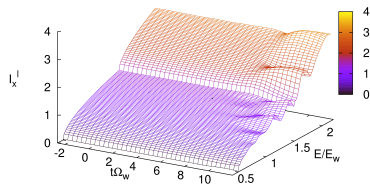


$$\beta = 1 \times 10^{-4} \text{ nm}^{-2}, \quad \beta = 4 \times 10^{-4} \text{ nm}^{-2}$$

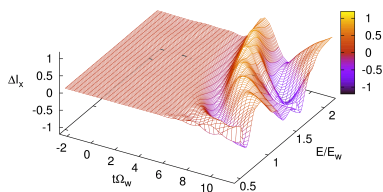
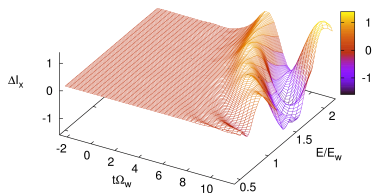
## Static conductance



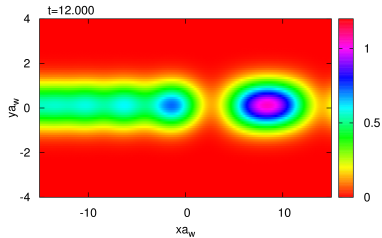
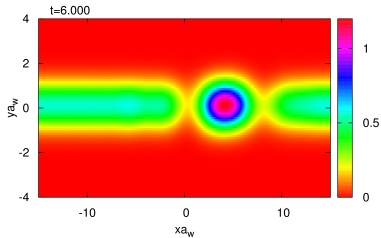
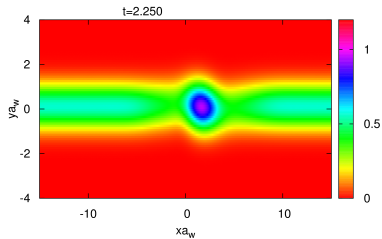
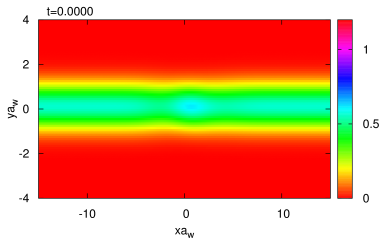
$I_x^l$  and  $I_x^r$ ,  $B = 0.1$  T,  $V_0 = -1$  meV



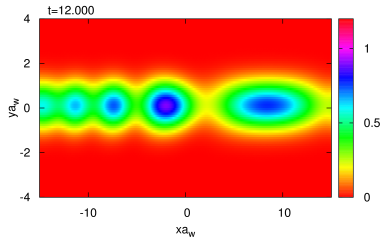
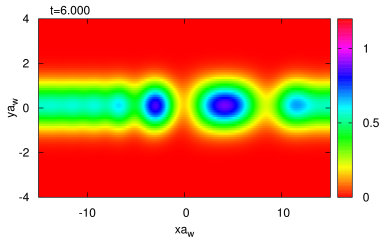
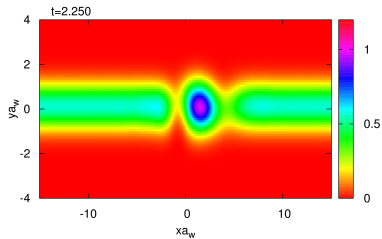
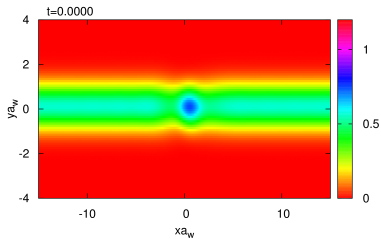
$I_x^l - I_x^r$ ,  $\beta = 1 \times 10^{-4} \text{ nm}^{-2}$ ,  $\beta = 4 \times 10^{-4} \text{ nm}^{-2}$



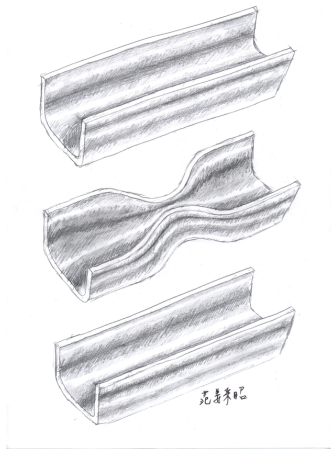
$$|\Psi|^2, B = 0.1 \text{ T}, V_0 = -1 \text{ meV}, \beta = 1 \times 10^{-4} \text{ nm}^{-2}, E = 0.75 E_w$$



$$|\Psi|^2, B = 0.1 \text{ T}, V_0 = -1 \text{ meV}, \beta = 4 \times 10^{-4} \text{ nm}^{-2}, E = 0.75 E_w$$

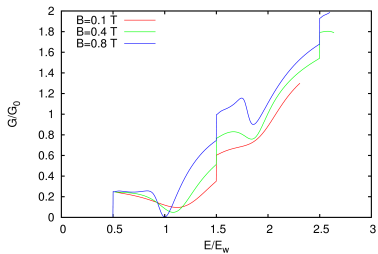
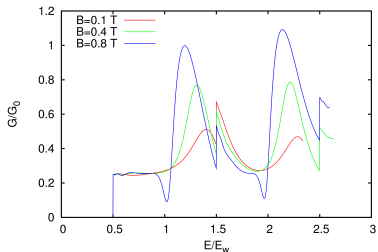


## Smooth hill-like pulse



$$\beta = 1 \times 10^{-4} \text{ nm}^{-2}, \quad \beta = 4 \times 10^{-4} \text{ nm}^{-2}$$

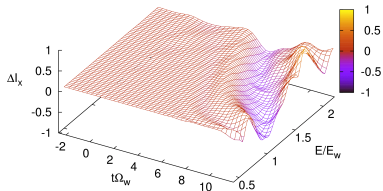
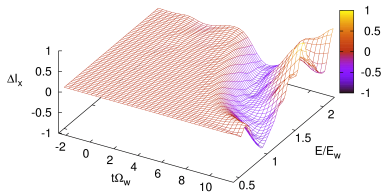
## Static conductance





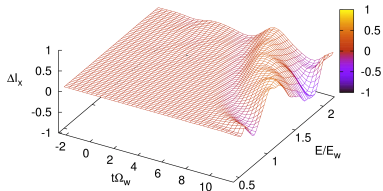
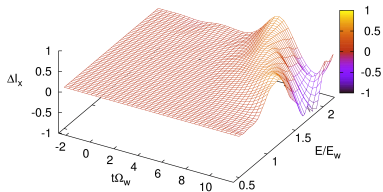
$\Delta I_x$ ,  $B = 0.4 \text{ T}$ ,  $V_0 = +1 \text{ meV}$ ,

$\beta = 1 \times 10^{-4} \text{ nm}^{-2}$ ,  $\beta = 4 \times 10^{-4} \text{ nm}^{-2}$

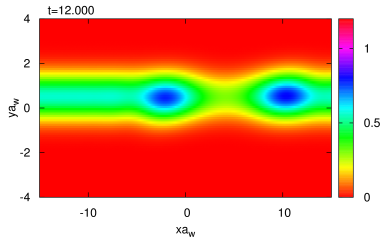
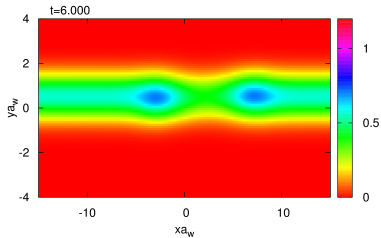
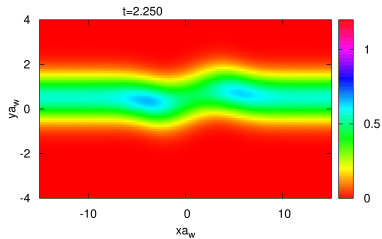
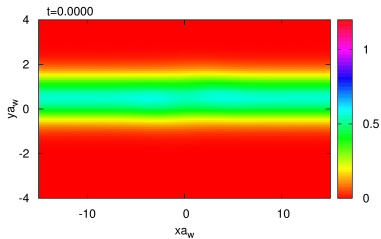


$\Delta I_x$ ,  $B = 0.4 \text{ T}$ ,  $V_0 = -1 \text{ meV}$ ,

$\beta = 1 \times 10^{-4} \text{ nm}^{-2}$ ,  $\beta = 4 \times 10^{-4} \text{ nm}^{-2}$



$$|\Psi|^2, B = 0.6 \text{ T}, V_0 = +1 \text{ meV}, \beta = 1 \times 10^{-4} \text{ nm}^{-2}, E = 0.63 E_w$$



# Conclusions

## Wave packet propagation

- Lifetimes of resonances
- Spreading of wave packets by resonances and magnetic field
- Mode-mixing  $\rightarrow$  onset of skipping orbits

## Current modulation

- Is modulation of current possible?
- Inelastic scattering
- Releasing of quasi-bound states?

- Magnetotransport in general smooth geometries
- Coulomb interaction, many-body effects?