

QEDFT = QED + DFT applied to an array of quantum dots in a photon cavity

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Experimental impetus...

physics

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Collective non-perturbative coupling of 2D electrons with high-quality-factor terahertz cavity photons

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The collective interaction of electrons with light in a highquality-factor cavity is expected to reveal new quantum phenomena¹⁻⁷ and find applications in quantum-enabled technologies^{8,9}. However, combining a long electronic coherence time, a large dipole moment, and a high quality-factor has proved difficult¹⁰⁻¹³. Here, we achieved these conditions simultaneously in a two-dimensional electron gas in a high-quality-factor terahertz cavity in a magnetic field. The vacuum Rabi splitting of cyclotron resonance exhibited a square-root dependence on the electron density, evidencing collective interaction. This splitting extended even where the detuning is larger than the resonance frequency. Furthermore, we observed a peak shift due to the normally negligible diamagnetic term in the Hamiltonian, Finally, the high-quality-factor cavity suppressed superradiant cyclotron resonance decay, revealing a narrow intrinsic linewidth of 5.6 GHz. High-quality-factor terahertz cavities will enable new experiments bridging the traditional disciplines of condensed-matter physics and cavity-based quantum optics.

nonresonant matter decay rate, which is usually the decoherence rate in the case of solids. Strong coupling is achieved when the splitting, 2g, is much larger than the linewidth, $(\kappa + \gamma)/2$, and ultrastrong coupling is achieved when g becomes a considerable fraction of ω_0 . The two standard figures of merit to measure the coupling strength are $C \equiv 4g^2/(\kappa\gamma)$ and g/ω_0 , here, C is called the cooperativity parameter", which is also the determining factor for the onset of optical bistability through resonant absorption saturation³⁷. To maximize C and g/ω_0 , one should construct a cavity QED set- up that combines a large dipole moment (that is, large g), a small decoherence rate (that is, small γ), a large cavity Q factor (that is, small γ) and a small resonance frequency ω_0 .

Group III-V semiconductor quantum wells (QWs) provide one of the cleanest and most tunable solid-state environments with quantum-designable optical properties. Microcavity QW-excitonpolaritons represent a landmark realization of a strongly coupled light-condensed-matter system that exhibits a rich variety of coherent many-body phenomena³¹. However, the large values of ω_0 and clatively small dipole moments for interband transitions make it

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- 2DEG in GaAs-AlGaAs heterostructure
- FIR photon cavity
- External magnetic field



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Exact diagonalization, one photon mode

- $\hbar\omega = 0.8 \text{ meV}$
- 2 electrons, first photon replica



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Large electron system – 2DEG





No exact diagonalization possible

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- $\blacksquare \ \mathsf{QED} + \mathsf{DFT} = \mathsf{QEDFT}$
- Use and adapt functional: $E_{\rm xc}^{\rm GA}[n_e, \nabla n_e]$, proposed by Johannes Flick, Simple Exchange-Correlation Energy Functionals for Strongly Coupled Light-Matter Systems based on the Fluctuation-Dissipation Theorem (2021), arXiv:2104.06980 [physics.chem-ph]

Orbital magnetization is sensitive to charge polarizability

Test for effects on orbital magnetization, M_o , of a 2DEG in a quantum dot array \leftrightarrow ground state property

$$\begin{split} M_o + M_s = & \frac{1}{2c\mathcal{A}} \int_{\mathcal{A}} d\boldsymbol{r} \ \left(\mathbf{r} \times \mathbf{j}(\mathbf{r}) \right) \cdot \hat{\mathbf{e}}_z \\ & - \frac{g^* \mu_B^*}{\mathcal{A}} \int_{\mathcal{A}} d\boldsymbol{r} \ \sigma_z(\mathbf{r}) \end{split}$$

- EM-field randomly polarized in the 2DEG plane
- External magnetic field, $\boldsymbol{B} \neq 0$

$${f A}=L^2$$
, $L=100~{
m nm}$

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Model and EM functional

$$H = H_0 + H_{\text{Zee}} + V_{\text{H}} + V_{\text{per}} + V_{\text{xc}} + V_{\text{xc}}^{\text{EM}}$$

$$\begin{split} E_{\rm xc}^{\rm GA}[n_e, \boldsymbol{\nabla} n_e] &= \frac{1}{16\pi} \sum_{\alpha=1}^{N_p} |\lambda_{\alpha}|^2 \int d\boldsymbol{r} \frac{\hbar \omega_p(\boldsymbol{r})}{\sqrt{(\hbar \omega_p(\boldsymbol{r}))^2/3 + (\hbar \omega_g(\boldsymbol{r}))^2} + \hbar \omega_{\alpha}} \\ & (\hbar \omega_g)^2 = C \left| \frac{\boldsymbol{\nabla} n_e}{n_e} \right|^4 \frac{\hbar^2}{m^{*2}} \\ & (\hbar \omega_p(q))^2 = (\hbar \omega_c)^2 + \frac{2\pi n_e^2}{m^* \kappa} q + \frac{3}{4} v_{\rm F}^2 q^2 \\ & \omega_c = \left(\frac{eB}{m^* c}\right), \quad l^2 = \left(\frac{hc}{eB}\right) \end{split}$$

Select $N_p = 1$, $\hbar\omega_{\alpha} = 1.0$ meV, L = 100 nm, $m^* = 0.067m_e$, $\kappa = 12.4$, $g^* = 0.44$, and $q \approx k_{\rm F}/6 \approx |\nabla n_e|/n_e$. $\lambda_{\alpha}l$ is measured in meV^{1/2}

Commensurability

L and l are competing length scales - Hofstadter problem (Phys. Rev. B 14, 2239 (1976))

Magnetic flux through unit cell: $B\mathcal{A} = pq\Phi_0$, $\Phi_0 = hc/e$, $p, q \in \mathbf{N}$

$$\begin{array}{l} N_e = 2, \ pq = 1 & \to \\ \lambda_\alpha l = 0.050 \ {\rm meV}^{1/2} \\ \mu = -8.954 \ {\rm meV} \\ T = 1.0 \ {\rm K} \\ \hbar \omega_\alpha = 1.0 \ {\rm meV} \\ E_{\rm Zee} = 1.053 \times 10^{-2} \ {\rm meV} \end{array}$$





Polaritons emerge, pq = 1



900

$$V_{\rm xc}$$
, $V_{\rm xc}^{\rm EM}$, $[n_e(\lambda_{\alpha}) - n_e(0)]$, $pq = 4$, $\lambda_{\alpha} l = 0.050 \ {\rm meV}^{1/2}$



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Total energy

$$pq = 1, 4$$

0.1

0.1



Orbital magnetization, $M_0 = \mu_{\rm B}^*/L^2$, $\lambda_{\alpha} l = 0 \rightarrow 0.1 \text{ meV}^{1/2}$



500

Cavity-photon influence on orbital magnetization pq = 1,3 pq = 2,4



900

Summary

- QEDFT (GGA), 2DEG
- Electron polarizability
- External magnetic field
- Orbital magnetization, total energy
- Cavity-photon, bandstructure and lattice effects
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