

1. project

①

Consider a harmonic oscillator in 1D, with

$$H_0 = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2, \quad [x, p] = i\hbar$$

introduce lowering and raising operators, such that

$$x = \frac{a}{\sqrt{2}} (a_+ + a_-), \quad p = \frac{i\hbar}{\sqrt{2}a} (a_+ - a_-)$$

$$a_- |n\rangle = \sqrt{n} |n-1\rangle, \quad a_+ |n\rangle = \sqrt{n+1} |n+1\rangle$$

on the eigenstates of H_0 $|n\rangle$ with spectrum $E_n^0 = \hbar\omega(n + \frac{1}{2})$

$a = \sqrt{\frac{\hbar}{m\omega}}$ is the natural length scale.

$$H_0 = \hbar\omega_0 \left\{ a_+ a_- + \frac{1}{2} \right\}$$

$$H_0 |n\rangle = E_n^0 |n\rangle, \quad E_n^0 = \hbar\omega_0 \left(n + \frac{1}{2} \right), \quad n=0, 1, 2, \dots$$

We can find the matrix element

$$\langle n | \frac{x}{a} | m \rangle = \frac{1}{2} \sqrt{n+m+1} S_{|n-m|, 1}$$

So, in the basis $\{|n\rangle\}$ the operator x can be represented by an infinite matrix. In the same way H_0 is represented by an infinite diagonal matrix

$$\langle n | H_0 | m \rangle = \hbar\omega_0 \left(n + \frac{1}{2} \right) \delta_{n,m}$$

How about x^2 ?

$$\langle n | \left(\frac{x}{a}\right)^2 | m \rangle = \langle n | \left(\frac{x}{a}\right) \left(\frac{x}{a}\right) | m \rangle$$

Now use the completeness relation

$$\mathbb{1} = \sum_{p=0}^{\infty} |p\rangle \langle p|$$

$$\langle n | \left(\frac{x}{a}\right)^2 | m \rangle = \sum_{p=0}^{\infty} \langle n | \left(\frac{x}{a}\right) | p \rangle \langle p | \left(\frac{x}{a}\right) | m \rangle$$

If we had noted the matrix for the operator by $\#$ then we see from this that the matrix for x^2 is just $\# \cdot \#$ where matrix multiplication is used.

a) find the energy spectrum for

$$H = H_0 + \lambda \hbar \omega \left(\frac{x}{a}\right)^4$$

for $\lambda \in [0, 1.2]$ for an appropriately truncated basis. test the truncation.

b) The eigenstates $|\alpha\rangle$ of H are a linear combination of the original basis

$$|\alpha\rangle = \sum_{n=0}^{\infty} C_{\alpha n} |n\rangle$$

and the coefficients $C_{\alpha n}$ come from the eigenvectors of H . Show with column-graphs the contribution of basis states to the 3 lowest energy states $|\alpha\rangle$. $|C_{\alpha n}|^2$ v.s. n

c) More complex functions of X

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There is a unitary transformation U that can get X on a diagonal form

$$X = U X_d U^\dagger = U \text{diag}\{\lambda_1, \dots, \lambda_n\} U^\dagger$$

eigenvalues of X

$$\rightarrow f(X) = U \text{diag}\{f(\lambda_1), \dots, f(\lambda_n)\} U^\dagger$$

where λ_i is the i^{th} eigenvalue of X and U is the unitary matrix formed of its eigenvectors.

We can scale

$$H_0 = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

using \hbar , and for fun we can also use $p = \hbar k$ such that the dimension of k : $[k] \sim \frac{1}{L}$

$$H_0 = \frac{(\hbar k a)^2}{2 m a^2} + \frac{a^2}{2} m \omega^2 \left(\frac{x^2}{a^2} \right) = \hbar \omega \left\{ \frac{(k a)^2}{2} + \frac{1}{2} \left(\frac{x^2}{a} \right) \right\} \quad (6)$$

The potential of the harmonic oscillator is thus

$$V_0(x) = \frac{1}{2} \hbar \omega \left(\frac{x}{a} \right)^2$$

Can we use the method to find the spectrum for

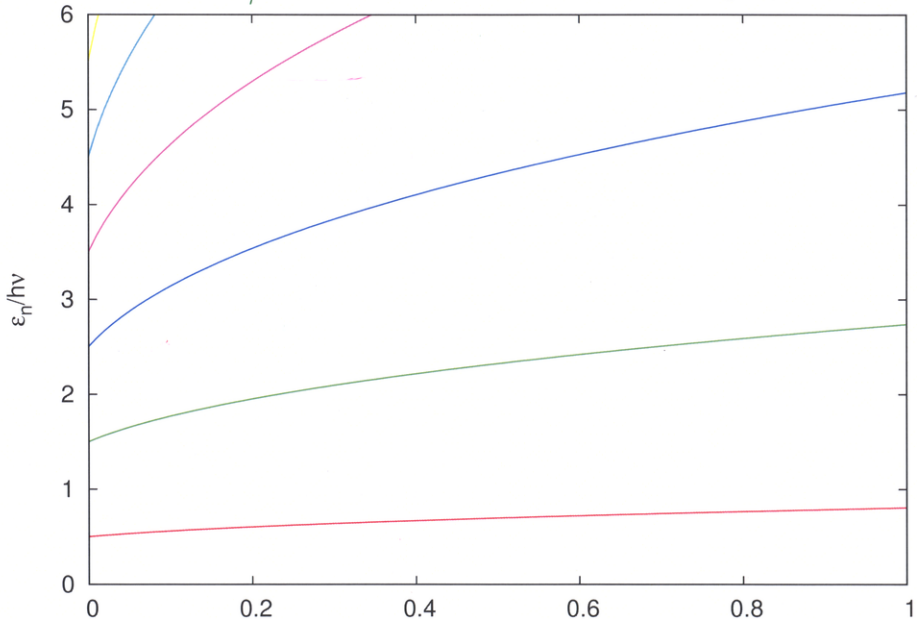
$$V(x) = \hbar \omega \left[\frac{1}{2} \left(\frac{x}{a} \right) \cdot \tanh \left(\frac{x}{a} \right) - \frac{1}{2} \left(\frac{x}{a} \right)^2 \right]$$

How large basis do we need to get a reasonable accuracy for the 10 lowest states in $V(x)$?

Check $V(x)$ in graph, notice that by an appropriate coefficient in the argument of \tanh we are really changing the parabolic potential into almost a V-shaped one.

Orbitöröf $H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2 + \lambda \tan\left(\frac{x}{a}\right)^4$

(2f)



Funkcio i granni 128 sęigirfalla $H_0 = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2$