

*Time-dependent transport through a nanostructure
via the generalized master equation, IV*

Viðar Guðmundsson

Science Institute, University of Iceland

`vidar@raunvis.hi.is`

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Coulomb interaction

$$\begin{aligned} \dot{\rho}_S(t) = & \frac{i}{\hbar} [H_S, \rho_S(t)] \\ & - \frac{1}{\hbar^2} \text{Tr}_{\text{res}} \left\{ \left[H_T(t'), \int_0^t dt' \left[U_0(t-t') H_T(t') U_0^\dagger(t-t') \right. \right. \right. \\ & \left. \left. \left. , e^{-i(t-t')H_S/\hbar} \rho_S(t') e^{+i(t-t')H_S/\hbar} \right] \right] \right\} \end{aligned}$$

can be rewritten as

$$\dot{\rho}_S(t) = \frac{i}{\hbar} [H_S, \rho_S(t)] - \frac{\chi(t)}{\hbar^2} \int_0^t dt' \chi(t') \sum_{l=\{L,R\}} \int dq K[t, t', q, l; \rho_S(t')]$$

with the kernel given by

$$\begin{aligned} K[t, t', q, l; \rho_S(t')] &= \left\{ [T^l(q), U_S^+(t-t') \mathcal{T}^{*l}(q) \rho_S(t') U_S(t-t')] e^{-i(t-t')\epsilon^l(q)/\hbar} \right. \\ &+ [U_S^+(t-t') \rho_S(t') \mathcal{T}^l(q) U_S(t-t'), \mathcal{T}^{*l}(q)] e^{+i(t-t')\epsilon^l(q)/\hbar} \left. \right\} (1 - f(\epsilon^l(q))) \\ &+ \left\{ [U_S^+(t-t') \rho_S(t') \mathcal{T}^{*l}(q) U_S(t-t'), \mathcal{T}^l(q)] e^{-i(t-t')\epsilon^l(q)/\hbar} \right. \\ &+ [\mathcal{T}^{*l}(q), U_S^+(t-t') \mathcal{T}^l(q) \rho_S(t') U_S(t-t')] e^{+i(t-t')\epsilon^l(q)/\hbar} \left. \right\} f(\epsilon^l(q)) \end{aligned}$$

With the Coulomb interaction we had

$$H_S = \sum_a E_a d_a d_a^\dagger + \frac{1}{2} \sum_{abcd} (ab|V_{\text{Coul}}|cd) d_a^\dagger d_b^\dagger d_d d_c$$

Diagonalization brings

$$|\mu\rangle = \mathcal{V}|\mu\rangle, \quad \mathcal{V}^\dagger|\mu\rangle = |\mu\rangle$$

and with

$$\tilde{\mathcal{T}}^l(q) = \mathcal{V}^\dagger \mathcal{T}^l(q) \mathcal{V}, \quad (\tilde{\mathcal{T}}^l(q))^* = \mathcal{V}^\dagger (\mathcal{T}^l(q))^* \mathcal{V}$$

we can transform the GME to the new interacting basis in which $\tilde{H}_S = \mathcal{V}^\dagger H_S \mathcal{V}$ is diagonal

Diagonalize $H_S \rightarrow$ **transform** GME unitarily \rightarrow **truncate** ρ_S and $\{|\mu\rangle\}$,
truncation is **not** possible before transformation

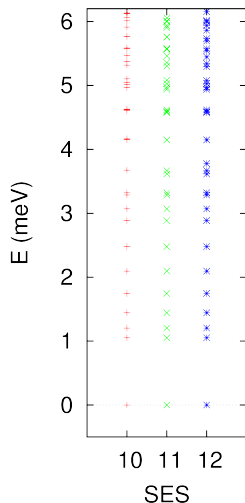
Operators can be represented in either basis and transformed between them

$$\langle Q_S(t) \rangle_I = \text{Tr}_S \{ \tilde{\rho}_S(t) \tilde{Q}_S \} = \text{Tr}_S \{ \tilde{\rho}_S(t) \mathcal{V}^\dagger Q_S \mathcal{V} \}$$

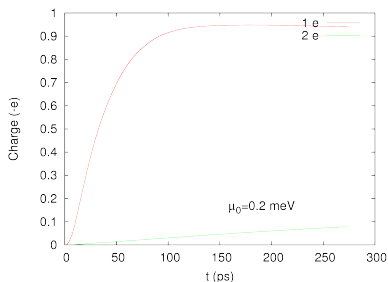
Finite quantum wire

Many-electron spectra

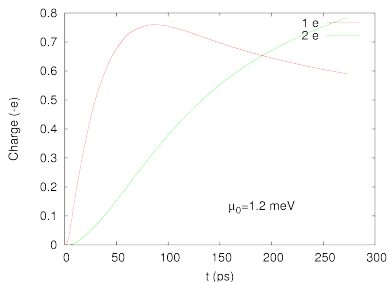
- $L_x = 300$ nm
- Parabolic confinement in y -direction, $\hbar\Omega_0 = 1.0$ meV
- Hard walls at $x = \pm L_x/2$
- $B = 1.0$ T
- GaAs parameters



$$\Delta\mu = \mu_L - \mu_R = 0.2 \text{ meV}$$



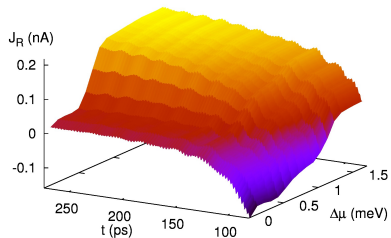
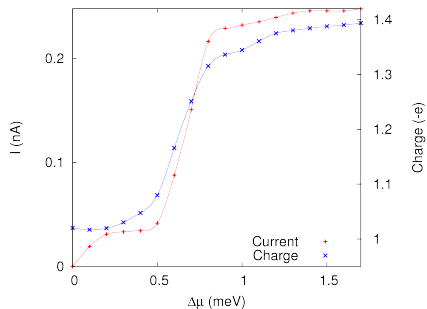
$$\Delta\mu = \mu_L - \mu_R = 1.2 \text{ meV}$$



Finite parabolic wire, Total charge

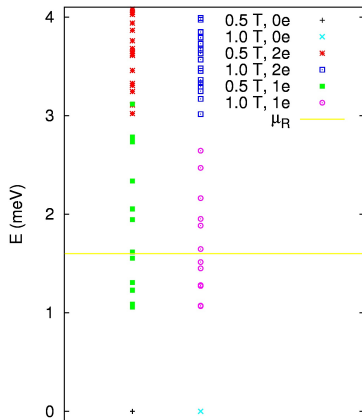
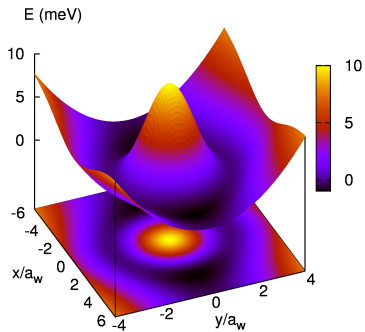
$$B = 1.0 \text{ T}, L_x = 300 \text{ nm}, \hbar\Omega_0 = 1.0 \text{ meV}$$

Total charge and current

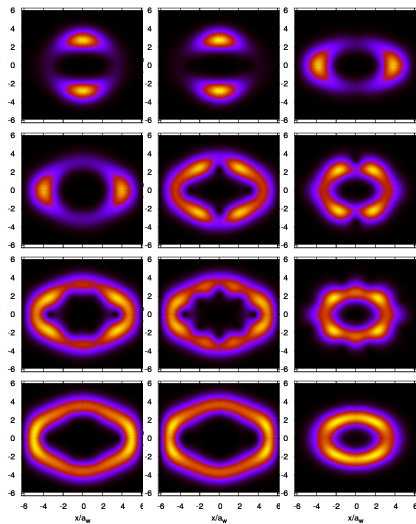


Coulomb blocking

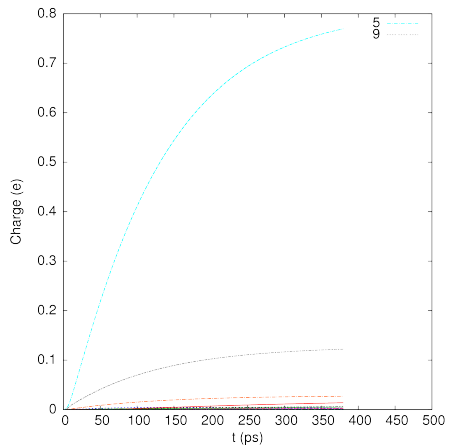
Embedded ring



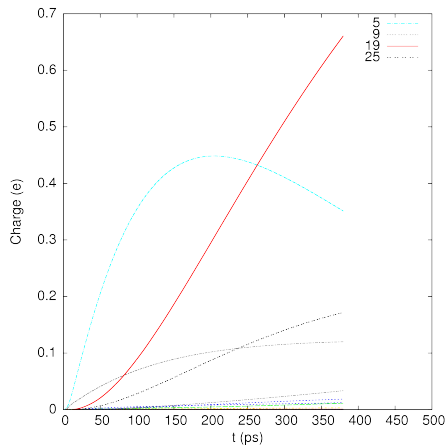
SES-probabilities



Coulomb interaction enhances occupation!

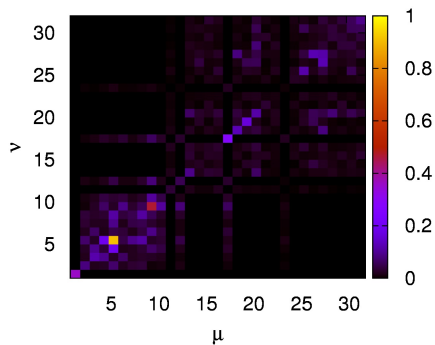


No interaction

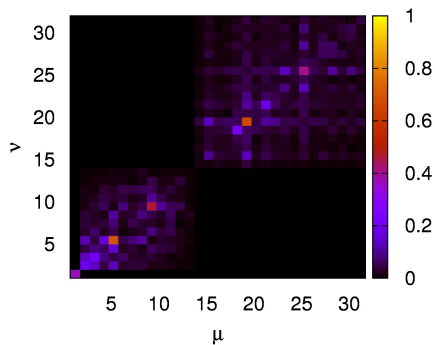


Interaction

Correlation enhanced by Coulomb interaction

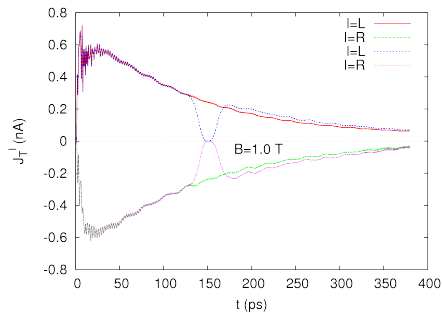
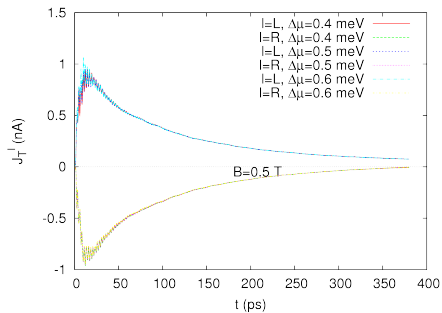


No interaction

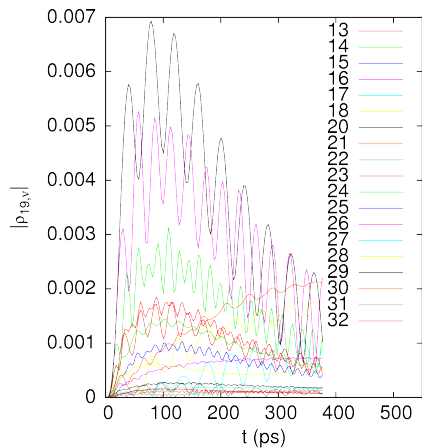


Interaction

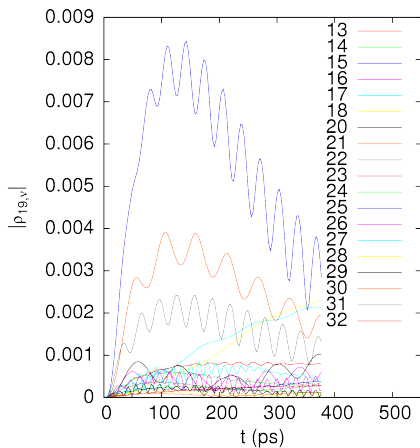
Current oscillations



Correlation oscillations

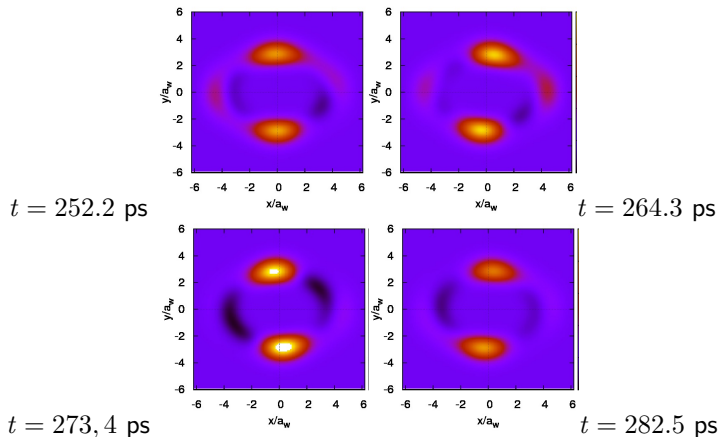


$B = 0.5$ T



$B = 1.0$ T

Correlation oscillations



$$n(\mathbf{r}, t) - n(\mathbf{r}, t - \delta t)$$
$$B = 1.0 \text{ T}$$

Summary

- GME-formalism
 - Bias
 - Weak coupling
 - Magnetic field
 - Many-electron formalism
 - General model
- Analytical + numerical
- Time-evolution, transients, steady state
- Coulomb interaction
 - Exact diagonalization
 - Coulomb blocking
 - Interaction enhances correlations
 - Correlation oscillations
- Geometry matters
- System extensions possible

Step-wise introduction of complexity – step-wise truncation of Fock spaces

References

- Transient charging: VM, AM, CST, VG, Phys. Rev. B81, 155442 (2010)
- Experimental systems: VM, AM, VG, Phys. Rev. B80, 205325 (2009)
- Correlations effects: VG, CST, OJ, VM, AM, Phys. Rev. B81, 205319 (2010)
- Dynamic cross correlation: VM, AM, VG, Phys. Rev. B82, 085311 (2010),