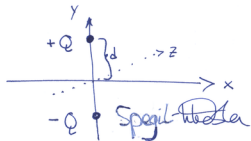


①

Punktladung  $Q$ a) finna  $\rho_s$ 

úr bók (4-37)

$$V(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{R_+} - \frac{1}{R_-} \right)$$

$$R_+ = \sqrt{x^2 + (y-d)^2 + z^2}$$

$$R_- = \sqrt{x^2 + (y+d)^2 + z^2}$$

$$\frac{\partial}{\partial y} \frac{1}{R_+} = -\frac{1}{2} \frac{2(y-d)}{(x^2 + (y-d)^2 + z^2)^{3/2}} \quad \text{①}$$

$$\frac{\partial}{\partial y} \frac{1}{R_-} = -\frac{1}{2} \frac{2(y+d)}{(x^2 + (y+d)^2 + z^2)^{3/2}}$$

I shall need the electric field at the plane  $y=0$ 

$$\vec{E}(x, y=0, z) = \left\{ -(\nabla V(\mathbf{r})) \cdot \hat{a}_y \right\} \hat{a}_y$$

$$= -\hat{a}_y \frac{Qd}{2\pi\epsilon_0 (x^2 + d^2 + z^2)^{3/2}}$$

and

$$\rho_s(x, 0, z) = \hat{a}_y \cdot \epsilon_0 \vec{E}(x, 0, z)$$

$$= -\frac{Qd}{2\pi (x^2 + d^2 + z^2)^{3/2}}$$

b) The total polarized charge on the surface  
In the plane polar coordinates are convenient

$$x^2 + z^2 = r^2$$

$$\rightarrow \rho_s(r) = -\frac{Qd}{2\pi(r^2+d^2)^{3/2}}$$

$$\int_0^{\infty} r dr \int_0^{2\pi} d\phi \rho_s(r) = -Qd \int_0^{\infty} \frac{r dr}{(r^2+d^2)^{3/2}} =$$

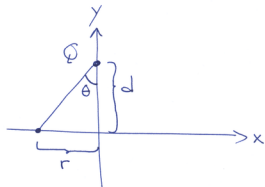
$$= +Qd \left. \frac{1}{\sqrt{r^2+d^2}} \right|_0^{\infty} = \underline{\underline{-Q}}$$

(2)

c) The force on  $\rho_s(r)$  due to  $Q$

(3)

$$\vec{F} = q \vec{E} \text{ generally, or } d\vec{F} = \vec{E} dq$$



$$d\vec{F} = \vec{E}(x,0,z) 2\pi r dr \rho_s(r)$$

due to  $q$ -Symmetry

x or z component = 0 of  $d\vec{F}$   
Cancel

Field of  $Q$  only

$$\begin{aligned} dF_y &= |d\vec{F}| \cdot \cos\theta = \underbrace{|E(x,0,z)|}_{\frac{Q}{4\pi\epsilon_0 R^2(y=0)}} \underbrace{\cos\theta}_{\frac{d}{\sqrt{r^2+d^2}}} 2\pi r dr \rho_s(r) \\ &= \frac{Q}{4\pi\epsilon_0 R^2(y=0)} \cdot 2\pi r dr \rho_s(r) \cdot \frac{d}{\sqrt{r^2+d^2}} \end{aligned}$$

$$F_y = \frac{(Qd)^2}{4\pi\epsilon_0} \int_0^\infty r dr \underbrace{\frac{1}{(r^2+d^2)}}_{\sim E} \underbrace{\frac{1}{\sqrt{r^2+d^2}}}_{\sim \cos\theta} \underbrace{\frac{1}{(r^2+d^2)^{3/2}}}_{\sim f_s}$$

$$= \frac{(Qd)^2}{4\pi\epsilon_0} \int r dr \frac{1}{(r^2+d^2)^3} = \frac{(Qd)^2}{4\pi\epsilon_0 4d^4}$$

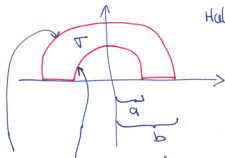
$$= \frac{Q^2}{4\pi\epsilon_0 (2d)^2}$$

We could not have used here  $\vec{E}$  found in a) since it includes the effects of  $Q$  and  $f_s$ , or  $Q$  and its image !!

which should answer item d)

(2)

(5)



Half-circular washer with thickness  $h$

Calculate  $R$  between the arched surfaces.

$\hookrightarrow$  only radial current

Choose potential  $V(a) = V_0, V(b) = 0$

$$\nabla^2 V(r) = 0 \iff \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) = 0$$

$$\rightarrow \frac{d}{dr} \left( r \frac{dV}{dr} \right) = 0 \rightarrow r \frac{dV}{dr} = C_1 \quad \text{or} \quad \frac{dV}{dr} = \frac{C_1}{r}$$

$$\rightarrow V(r) = C_1 \ln r + C_2$$

$\rightarrow$  B.C.

$$V_0 = C_1 \ln a + C_2$$

$$0 = C_1 \ln b + C_2$$

$$C_2 = - \frac{V_0 \ln(b)}{\ln(a/b)}$$

$$C_1 = \frac{V_0}{\ln(a/b)}$$

$$V(r) = V_0 \frac{\ln\left(\frac{r}{b}\right)}{\ln\left(\frac{a}{b}\right)} \quad \rightarrow \quad \vec{E}_r = -\hat{a}_r \frac{\partial V}{\partial r}$$

$$= -\hat{a}_r \frac{V_0}{\ln(a/b)}$$

$\vec{E}$  is in the direction of  $\hat{a}_r$   
 since  $\ln\left(\frac{a}{b}\right) < 0$  for  $a < b$

$$\vec{J}(r) = \nabla \vec{E}(r) \quad \rightarrow \quad I = \int_S \vec{J}(r) \cdot d\vec{s}$$

$$\rightarrow I = \int_0^{\pi} J(r) h r d\phi = - \frac{\pi \nabla h V_0}{\ln(a/b)}$$

$$R = \frac{V_0}{I} = - \frac{\ln\left(\frac{a}{b}\right)}{\pi \nabla h} = \frac{\ln\left(\frac{b}{a}\right)}{\pi \nabla h} > 0$$

(6)

3

7

$\mu_0$

$\uparrow \uparrow \bar{H}_0 = \hat{A}_z H_0$

← given external magnetic field



Generally

$$\bar{H} = \frac{1}{\mu} \bar{B}$$

$$\bar{B} = \mu \bar{H}$$

$$\bar{H} = \frac{\bar{B}}{\mu_0} - \bar{M}$$

$$\bar{B} = (\bar{H} + \bar{M}) \mu_0$$

$\mu_0$

$\uparrow \uparrow \bar{H}_0 = \hat{A}_z H_0$

no parallel field here and no surface current



B.C.

$$B_{1n} = B_{2n}$$

$$\mu_1 H_{1n} = \mu_2 H_{2n}$$

Use the B.C.

a) The slab has  $\mu$

$$\mu H_i = \mu_0 H_0 \rightarrow \bar{H}_i = \hat{a}_z H_i = \hat{a}_z \frac{\mu_0}{\mu} H_0$$

b) The slab has magnetization  $\bar{M}_i$

$$\rightarrow \bar{B}_i = (\bar{H}_i + \bar{M}_i) \mu_0 \quad \text{and outside} \quad \bar{B}_o = \mu_0 \bar{H}_o$$

B.C.

$$\bar{B}_o = \mu_0 H_0 = \bar{B}_i = (H_i + M_i) \mu_0$$

$$\rightarrow \bar{H}_i = \hat{a}_z (H_o - M_i)$$



④ Spherical wave

$$\vec{E} = \hat{a}_\theta \frac{E_0}{R} \sin \theta \cos(\omega t - kR)$$

In phasor notation

$$\vec{E} = \hat{a}_\theta \frac{E_0}{R} \sin \theta e^{-ikR} \quad \text{only } \theta\text{-Component}$$

a) find  $\vec{H}$ . Use Faradays Law

$$\frac{j}{\omega \mu_0} \nabla \times \vec{E} = \vec{H}$$

$$\nabla \times \vec{E} = \hat{a}_\phi \frac{1}{R} \frac{\partial}{\partial R} (RE_\theta) = \hat{a}_\phi (-ik) E_\theta$$

$$\rightarrow \vec{H} = \hat{a}_\phi \frac{k}{\omega \mu_0} E_\theta$$

⑨

b) 
$$S_{ave} = \frac{1}{2} \operatorname{Re} (\bar{\mathbf{E}} \times \bar{\mathbf{H}}^*)$$

$$= \frac{1}{2} \operatorname{Re} \left\{ \hat{\mathbf{a}}_{\theta} \times \hat{\mathbf{a}}_{\phi} \left( \frac{k}{\omega \mu_0} |E_{\theta}|^2 \right) \right\}$$

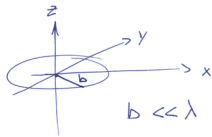
$$= \frac{1}{2} \hat{\mathbf{a}}_r \frac{k E_0^2}{\omega \mu_0 R^2} \sin^2 \theta$$

c) The total time-average power (flowing out)

$$P_{ave}^{out} = \int_0^{4\pi} R^2 d\Omega S_{ave} \cdot \hat{\mathbf{a}}_r = 2\pi \int_0^{\pi} R^2 \sin \theta d\theta \bar{S}_{ave} \cdot \hat{\mathbf{a}}_r$$
$$= 2\pi \frac{1}{2} \frac{R^2 k E_0^2}{\omega \mu_0 R^2} \frac{4}{3} = \frac{4\pi}{3} \frac{k E_0^2}{\omega \mu_0} = \frac{4\pi}{3} \frac{E_0^2}{\eta_0}$$

nonvanishing as  $R \rightarrow \infty$

(5)

Dipole  $I(t) = I_0 \cos(\omega t)$  $b \ll \lambda$  far field

$$E_{\phi} = \frac{\omega \mu_0 m}{4\pi R} \left( \frac{e^{-i\beta R}}{R} \right) \beta \sin\theta$$

$$H_{\theta} = -\frac{\omega \mu_0 m}{4\pi \eta_0} \left( \frac{e^{-i\beta R}}{R} \right) \beta \sin\theta$$

with

$$m = I_0 \pi b^2$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} \approx 377 \Omega$$

(11)

a)

$$P_r = \oint \overline{S}_{ave} \cdot d\vec{s} \quad , \quad \overline{S}_{ave} = \frac{1}{2} \text{Re} (\vec{E} \times \vec{H}^*)$$

$$P_r = -\frac{1}{2} \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta E_\phi H_\theta^* R^2$$

$$= \frac{1}{2} \left( \frac{\omega \mu_0 \mu}{4\pi} \right)^2 \frac{\beta^2}{\epsilon_0} 2\pi \int_0^{\pi} \sin^3\theta d\theta$$

$$= \frac{2}{3} \left( \frac{\omega \mu_0 \mu}{4\pi} \right) \frac{\beta^2}{\epsilon_0} 2\pi, \quad \beta = \frac{2\pi}{\lambda}$$

$$\begin{aligned}
 &= \frac{2}{3} \left( \frac{\omega \mu_0 I_0 \pi b^2}{4\pi} \right)^2 2\pi \left( \frac{2\pi}{\lambda} \right)^2 \frac{1}{60\pi} \\
 &= \frac{2}{3} \left( \frac{\omega \mu_0 I_0 \pi b^2}{2\lambda} \right)^2 \frac{1}{60} = \frac{1}{2} I_0^2 \left\{ \frac{\omega^2 \mu_0^2 \pi^2 b^4}{180 \lambda^2} \right\} \\
 &= \frac{1}{2} I_0^2 R_r \rightarrow R_r = \frac{\omega^2 \mu_0^2 \pi^2 b^4}{180 \lambda^2} = 80\pi^2 \left( \frac{\beta \pi b}{\lambda} \right)^2
 \end{aligned}$$

b)

$$\eta_r = \frac{R_r}{R_r + R_e}$$

loss in wire  
 we know  $V_0$ , length of wire  $2\pi b = l$  and cross section  $S = \pi a^2$

$$\rightarrow R_e = \frac{l}{\sigma_0 s} = \frac{2\pi b}{\sigma_0 \pi a^2} = \frac{2b}{a^2 \sigma_0}$$

only true for low frequency, for higher  $f$  we have surface conduction

$$R_s = \sqrt{\frac{\omega \mu_0}{2\sigma_0}} \quad \text{see (11-49)}$$

$$R_e = R_s \frac{2\pi b}{2\pi a} = R_s \frac{b}{a}$$

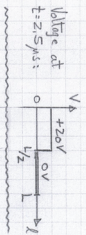
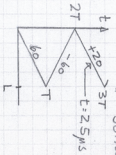
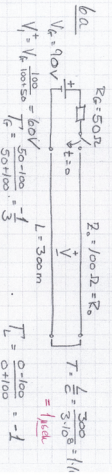
$$\rightarrow Y_r = \frac{1}{1 + \frac{R_e}{R_r}} = \frac{1}{1 + \frac{b \sqrt{\frac{\omega \mu_0}{2\sigma_0}}}{80\pi^2 \left(\frac{B\pi b}{\lambda}\right)^2}}$$

c) If  $b$  and  $l$  are fixed it would be ~~best~~ to have several windings of wire in the antenna  $\leftarrow$  see  $R_r$

d) We have fixed  $I_0 \cos(\omega t)$  in the beginning  
 $\rightarrow R_r$ , or higher value of  $R_r$   
represent how large resistance we would need to dissipate the radiated power

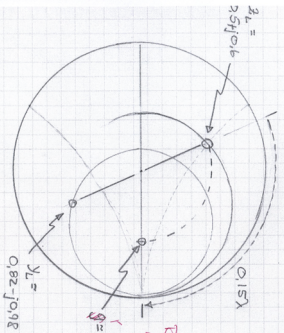
SOLUTIONS FOR TRANSMISSION LINE QUESTIONS

6a



6b

$Z_L = 50 + j60 \Omega \Rightarrow Z_L = \frac{1}{R_0} \cdot Z_L = 0.5 + j0.6$



$d = 0.15 \lambda$

$Y_L = \frac{1}{R_0} Y_L = 8.2 - j0.8 mS$

$R \approx R_0 \cdot \rho \approx 290 \Omega$

*Varloedmitt i afslutning*