

4.17

Í stað málisortu

$$V(r) = -\frac{e^2}{4\pi\epsilon_0 r}$$

skodum við

$$V(r) = -\frac{GMm}{r}$$

M: massi sólar

m: massi jarðar

p.a. í jöfunum fyrir

H-atómið þarfjum við

að setja

$$GMm \leftarrow \frac{e^2}{4\pi\epsilon_0}$$

b) fyrir H-atóm er Bohr geislin ①

$$a_H = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \left(\frac{4\pi\epsilon_0}{e^2}\right) \frac{\hbar^2}{m_e}$$

fyrir J-S kerfið fast þá

$$a_G = \frac{\hbar^2}{GMm^2}, \quad G = 6.673 \cdot 10^{-11} \frac{\text{m}^3}{\text{kg s}^2}$$

$$= \frac{(1.055 \cdot 10^{-34} \text{ J s})^2}{6.673 \cdot 10^{-11} \frac{\text{m}^3}{\text{kg s}^2} \cdot 2 \cdot 10^{30} \text{ kg} \cdot (6 \cdot 10^{24} \text{ kg})^2}$$

$$= 2.32 \cdot 10^{-138} \text{ m}$$

c) fyrir H-atóm var

$$E_n = -R_y^H \cdot \frac{1}{n^2}$$

$$R_y^H = \left[ \frac{m_e}{2h^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \right]$$

$$\rightarrow R_y^G = \frac{m}{2h^2} (GMm)^2$$

$$E_n^G = -R_y^G \frac{1}{n^2} = -\frac{m}{2h^2} (GMm)^2 \frac{1}{n^2}$$

Orka þá er á hringsreytingu (2)

$$E = \frac{1}{2} m v^2 - \frac{GMm}{r_0} = \text{fasti}$$

notum  $F = ma$  t.p.a. fuma  $v$

$$m \left( \frac{v^2}{r_0} \right) = \frac{GMm}{r_0^2} \rightarrow v^2 = \frac{GM}{r_0}$$

$$\rightarrow E = \frac{GMm}{2r_0} - \frac{GMm}{r_0}$$

$$= -\frac{GMm}{2r_0}$$

setjum jafnt

$$E_n^G = E$$

$$\frac{m}{2h^2} (GMm)^2 \frac{1}{n^2} = \frac{GMm}{2r_0}$$

$$\rightarrow \frac{m}{2\hbar^2} G M m \frac{1}{n^2} = \frac{1}{2r_0}$$

$$\rightarrow n^2 = \frac{r_0 G M m^2}{\hbar^2} = \frac{r_0}{a_G}$$

$$\rightarrow n = \sqrt{\frac{r_0}{a_G}}$$

$$\rightarrow n \approx \sqrt{\frac{150 \cdot 10^9 \text{ m}}{2.32 \cdot 10^{-138} \text{ m}}}$$

$$\approx 2.5 \cdot 10^{74}$$

d) Jördin geistlar  $n \rightarrow n-1$   
hve nutil orbita losnar

$$\Delta E^G = |E_{n-1}^G - E_n^G|$$
$$= + R_y^G \left\{ \frac{1}{(n-1)^2} - \frac{1}{n^2} \right\}$$
$$= + R_y^G \frac{1}{n^2} \left\{ \frac{n^2}{(n-1)^2} - 1 \right\}$$

$$= |E_n| \cdot \left\{ \frac{1}{(1-\frac{1}{n})^2} - 1 \right\}$$

$$\approx |E_n| \cdot \left\{ 1 + \frac{2}{n} - 1 \right\}$$

$$= |E_n| \cdot \frac{2}{n}$$

$$\Delta E^G = 2 \cdot R_y^G \frac{1}{n^3}$$

(3)

$$R_Y^G = \frac{m}{2\hbar^2} (GMm)^2 = \frac{(6 \cdot 10^{24} \text{ kg})^3 (6.673 \cdot 10^{-11} \frac{\text{m}^3}{\text{kg s}^2})^2 (2 \cdot 10^{30} \text{ kg})^2}{2 \cdot (1.055 \cdot 10^{-34} \text{ Js})^2} \quad (4)$$

$$\approx 1.73 \cdot 10^{182} \text{ J}$$

$$\rightarrow \Delta E^G \approx 2 \cdot R_Y^G \frac{1}{h^3} \sim 2.1 \cdot 10^{-41} \text{ J.}$$

Hvæða bylgjulengd myndi þetta samsvara fyrir þyngðkreið?

$$\begin{aligned} \Delta E^G &= h\omega = h \frac{\omega}{2\pi} = h\nu \\ &= h \frac{c}{\lambda} \end{aligned}$$

$$\begin{aligned} \rightarrow \lambda &= \frac{hc}{\Delta E^G} = \frac{(1.055 \cdot 10^{-34} \text{ Js} \cdot 2\pi) 3 \cdot 10^8 \text{ m}}{2.1 \cdot 10^{-41} \text{ J}} \approx 9.48 \cdot 10^{15} \text{ m} \\ &\approx 1 \text{ ljósár} \end{aligned}$$

4.55

Rafjandi vetni  $L$ 

$$R_{21} \left\{ \sqrt{\frac{1}{3}} Y_{10} X_+ + \sqrt{\frac{2}{3}} Y_{11} X_- \right\}$$

a) Málindarstöður  $L^2$ ?

Báðir þessar ástandsinns eru eiginástand  $L^2$  með eigingildi  $\hbar^2 l(l+1) = 2\hbar^2$ , líkúndin á þeim málindast. eru því 1.

b) Málindarstöður  $L_z$ ?

Báðir þessar ástandsinns eru eiginástand  $L_z$ , en með mismunandi eigingildi. Við fáum því

$\hbar \cdot 0$  með líkúndum  $\frac{1}{3}$

$\hbar$  með líkúndum  $\frac{2}{3}$

heildarlíkur eru  
því 1

c)  $S^2$  ?

Bádir spuna þetta ástandið er eiginástand  $S^2$  með sama eigingildit  $\rightarrow$  maling gefur

$$\hbar^2 \frac{1}{2} \left( \frac{1}{2} + 1 \right) = \frac{3\hbar^2}{4} \text{ með líkum } 1$$

d)  $S_z$

Bádir þetta er eiginástand  $S_z$ , en með mismunandi eigingildum, þú fóst

$$\begin{aligned}
 &+ \frac{\hbar}{2} \text{ með líkum } \frac{1}{3} \\
 &- \frac{\hbar}{2} \text{ með líkum } \frac{2}{3}
 \end{aligned}$$

e)  $J^2$

Hér vandast málið, ástandin eru ekki eigin ástand  $J^2$ .  
Við þurfum því að tala ástandið  $i$  þeim

$$\left. \begin{array}{l} Y_{10} \chi_+ \text{ stendur fyrir } |1,0\rangle |\frac{1}{2}, \frac{1}{2}\rangle \\ Y_{11} \chi_- \text{ ——— } |1,1\rangle |\frac{1}{2}, -\frac{1}{2}\rangle \end{array} \right\} \begin{array}{l} \text{tökum} \\ |l, m\rangle |S, m_s\rangle \end{array}$$

Í báðum tilfellum fast að  $m + m_s = \frac{1}{2}$ , því  $i$  samanturdi við jöfnu (4.184) fast aðeins að  $J = \frac{3}{2}$  eða  $\frac{1}{2}$

Við þurfum því að nota (4.186) t.p. a finna

$$|l, m\rangle |S, m_s\rangle = \sum_j C_{m, m_s, m_j}^{l, S, j} |j, m_j\rangle$$

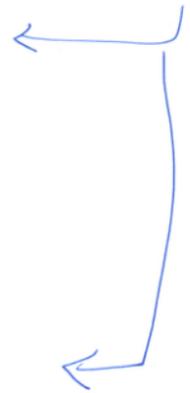
$$|1,0\rangle | \frac{1}{2}, \frac{1}{2} \rangle = \sum_j C_{0, \frac{1}{2}, m_j}^{1, \frac{1}{2}, j} |j, m_j\rangle$$

$$= \sqrt{\frac{2}{3}} | \frac{3}{2}, \frac{1}{2} \rangle - \sqrt{\frac{1}{3}} | \frac{1}{2}, \frac{1}{2} \rangle$$

Samkvæmt töflu 4.8

$$|1,1\rangle | \frac{1}{2}, -\frac{1}{2} \rangle = \sum_j C_{1, -\frac{1}{2}, m_j}^{1, \frac{1}{2}, j} |j, m_j\rangle$$

$$= \sqrt{\frac{1}{3}} | \frac{3}{2}, \frac{1}{2} \rangle + \sqrt{\frac{2}{3}} | \frac{1}{2}, \frac{1}{2} \rangle$$



og heitbar kvæðifunga hluti ástandsins er því í  $\{|j, m_j\rangle\}$ -grunn

$$\left[ \sqrt{\frac{1}{3}} |1,0\rangle | \frac{1}{2}, \frac{1}{2} \rangle + \sqrt{\frac{2}{3}} |1,1\rangle | \frac{1}{2}, -\frac{1}{2} \rangle \right]$$

$$= \sqrt{\frac{1}{3}} \left\{ \sqrt{\frac{2}{3}} | \frac{3}{2}, \frac{1}{2} \rangle - \sqrt{\frac{1}{3}} | \frac{1}{2}, \frac{1}{2} \rangle \right\} + \sqrt{\frac{2}{3}} \left\{ \sqrt{\frac{1}{3}} | \frac{3}{2}, \frac{1}{2} \rangle + \sqrt{\frac{2}{3}} | \frac{1}{2}, \frac{1}{2} \rangle \right\}$$

$$= \frac{2\sqrt{2}}{3} \overset{j, m_j}{|\frac{3}{2}, \frac{1}{2}\rangle} + \frac{1}{3} \overset{j, m_j}{|\frac{1}{2}, \frac{1}{2}\rangle}$$

Astandið er þú líklegt í eiginástandum  $J^2$  og  $J_z$  og við sjáum að mæling gefur

$$\text{gildi } \frac{3}{2} \left(\frac{5}{2}\right) \hbar^2 \quad \text{með líkum } \frac{4 \cdot 2}{9} = \frac{8}{9}$$

$$\frac{1}{2} \left(\frac{3}{2}\right) \hbar^2 \quad \text{með líkum } \frac{1}{9}$$

f) mæling á  $J_z$  gefur

$$\frac{\hbar}{2} \quad \text{með líkum } 1$$

Það er þó líklegt  
 að eiginástand  $J_z$   
 með sama eigin-  
 gildi

(9)

g) Ástandið var

$$R_{21} \left\{ \sqrt{\frac{1}{3}} Y_{10} \chi_+ + \sqrt{\frac{2}{3}} Y_{11} \chi_- \right\}$$

(10)

líkindaþéttleikinn

þegar þú ert farna eindina e

$$|R_{21}|^2 \left\{ \frac{1}{3} |Y_{10}(\vartheta)|^2 + \frac{2}{3} |Y_{11}(\vartheta)|^2 \right\}$$

þar sem  $\chi_+^* \chi_- = 0$  og  $\chi_-^* \chi_+ = 0$   
 $\chi_+ \chi_+ = 1$  og  $\chi_- \chi_- = 1$

$$\frac{1}{24} \frac{1}{a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}} \left\{ \frac{1}{4\pi} \cos^2 \theta + \frac{2}{8\pi} \sin^2 \theta \right\}$$

$$= \frac{1}{96\pi} \frac{1}{a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}}$$

w) liturnar  $\bar{a}$  og mola fjarlægð  $r$  og  $S_z$  með  $+\frac{\hbar}{2}$

(11)

$$|R_{21}|^2 \frac{1}{3} \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta |Y_{10}|^2$$

$$= \frac{1}{24} \frac{1}{a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}} \frac{2\pi}{3} \int_0^{\pi} d\theta \sin\theta \cos^2\theta \frac{3}{4\pi}$$

$$= \frac{1}{48} \frac{1}{a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}} \int_0^{\pi} d\theta \sin\theta \cos^2\theta$$

$$= \frac{1}{48} \frac{1}{a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}} \frac{2}{3} = \frac{1}{72a^3} \left(\frac{r}{a}\right)^2 e^{-\frac{r}{a}}$$