

⑥.1

①

Oändligt begränsade brynnar med samma energi

$$\psi_n^0(x) = \sqrt{\frac{2}{a}} \sin\left(n\pi \frac{x}{a}\right)$$

$$E_n^0 = E_1 \cdot n^2, \quad E_1 = \frac{\pi^2 \hbar^2}{2ma^2}$$

Kerfö ar truföe med  $H' = \alpha \delta(x - \frac{a}{2})$

a) Fyrsta Stigs truflem  $\bar{a}$  orku-rö-fime

$$E_n = E_n^0 + \langle n | V | n \rangle$$

$$\langle n | V | n \rangle = 2\frac{\alpha}{a} \int_0^a dx \sin^2\left(n\pi \frac{x}{a}\right) \delta\left(x - \frac{a}{2}\right)$$

$$\langle n | v | n \rangle = 2 \frac{\alpha}{a} \sin^2\left(\frac{n\pi}{2}\right) = \begin{cases} 0 & \text{ef } n = 2, 4, 6, \dots \\ \frac{2\alpha}{a} & \text{ef } n = 1, 3, 5, \dots \end{cases} \quad (2)$$

ástöndin með jöfnu  $n$  eru með nülstöð í  $x = \frac{a}{2}$ , miðjum brunn, eru andsamhverf um þennan punkt.

b) Finna lagðu löðina í löðrettingunni á grunnástandinu

$$|\psi_1'\rangle = \sum_{m=2}^{\infty} \frac{\langle \psi_m | H' | \psi_1^0 \rangle}{E_1^0 - E_m^0} |\psi_m^0\rangle$$

Þó  $m=1$  er sleppt, þurfum einu að sjá hvaða  $m$  er í reum summum yfir

partikel terikat

(3)

$$\langle \psi_m | H' | \psi_1^0 \rangle = 2 \frac{\alpha}{a} \int_0^a dx \sin(\pi \frac{x}{a}) \delta(x - \frac{a}{2}) \sin(m\pi \frac{x}{a})$$

$$= 2 \frac{\alpha}{a} \sin(\frac{\pi}{2}) \sin(m \frac{\pi}{2}) = \begin{cases} 2 \frac{\alpha}{a} \sin(\frac{m\pi}{2}) & \text{if } m=1,3,5,7 \\ 0 & \text{if } m=2,4,6 \end{cases}$$

Enn fremur er

$$E_1^0 - E_m^0 = E_1 - E_1 m^2 = E_1 (1 - m^2)$$

$$\rightarrow |\psi_1^1\rangle = \sum_{m=3,5,7,\dots}^{\infty} \left\{ \frac{2(\frac{\alpha}{a})}{E_1} \right\} \frac{\sin(\frac{m\pi}{2})}{1 - m^2}$$

Fyrsta 3 líðirnar eru

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$$|\Phi_1'\rangle \approx \frac{2\left(\frac{a}{a}\right)}{E_1} \left\{ \frac{-|\psi_3^0\rangle}{1-9} + \frac{|\psi_5^0\rangle}{1-25} + \frac{-|\psi_7^0\rangle}{1-49} \right\}$$

$$= \frac{2\left(\frac{a}{a}\right)}{E_1} \left\{ \frac{|\psi_3^0\rangle}{8} - \frac{|\psi_5^0\rangle}{24} + \frac{|\psi_7^0\rangle}{48} \right\}$$

$$\langle x | \Phi_1' \rangle \approx \frac{2\left(\frac{a}{a}\right)}{E_1} \left| \frac{2}{a} \right| \left\{ \frac{\sin\left(\frac{3\pi x}{a}\right)}{8} - \frac{\sin\left(\frac{5\pi x}{a}\right)}{24} + \frac{\sin\left(\frac{7\pi x}{a}\right)}{48} \right\}$$

Trúfnum er spegilsamhverft um punktin  $x = \frac{a}{2} \rightarrow$  Þessins bylgjufall sem eru með sömu samhverfu eru líka saman í trúföðu bylgjufallið

(6.6) Tvö "goda" ortofunktionärsständ eru  $|\psi_a^0\rangle$  og  $|\psi_b^0\rangle$  eru margföld

$$|\psi_{\pm}^0\rangle = \alpha_{\pm} |\psi_a^0\rangle + \beta_{\pm} |\psi_b^0\rangle \quad \left| \begin{array}{l} \langle \psi_a^0 | \psi_c^0 \rangle = \delta_{ac} \\ \text{med } c = a, b \end{array} \right.$$

þar sem  $\alpha_{\pm}$  og  $\beta_{\pm}$  eru ákvörðuð með jöfnu (6.22)

a) Sýna að  $\langle \psi_+^0 | \psi_-^0 \rangle = 0$

$$\left. \begin{array}{l} \text{ákvörðunargjöfnur eru þær} \\ \alpha_{\pm} W_{aa} + \beta_{\pm} W_{ab} = \alpha_{\pm} E'_{\pm} \\ \alpha_{\pm} W_{ba} + \beta_{\pm} W_{bb} = \beta_{\pm} E'_{\pm} \end{array} \right\} (*) \quad \begin{array}{l} \langle \psi_+^0 | \psi_-^0 \rangle \\ = \left\{ \alpha_+^* \langle \psi_a^0 | + \beta_+^* \langle \psi_b^0 | \right\} \\ \cdot \left\{ \alpha_- |\psi_a^0\rangle + \beta_- |\psi_b^0\rangle \right\} \\ = \alpha_+^* \alpha_- + \beta_+^* \beta_- \end{array}$$

Notum (\*)  $\downarrow$  p.a. unitar  $\beta_{\pm}$  i  $\alpha_{\pm}$

(6)

$$\beta_{\pm} W_{ab} = \alpha_{\pm} E_{\pm}' - \alpha_{\pm} W_{aa} \rightarrow \beta_{\pm} = \frac{\alpha_{\pm} (E_{\pm}' - W_{aa})}{W_{ab}}$$

$$\rightarrow \beta_{+}^{*} \beta_{-} = \frac{\alpha_{+}^{*} \alpha_{-}}{W_{ab}^{*} W_{ab}} (E_{+}' - W_{aa}^{*})(E_{-}' - W_{aa}), \quad W_{aa}^{*} = W_{aa}$$

pu fast

$$\begin{aligned} \langle 2\psi_{+}^{0} | 2\psi_{-}^{0} \rangle &= \frac{\alpha_{+}^{*} \alpha_{-}}{|W_{ab}|^2} \left[ |W_{ab}|^2 + (E_{+}' - W_{aa})(E_{-}' - W_{aa}) \right] \quad (***) \\ &= \frac{\alpha_{+}^{*} \alpha_{-}}{|W_{ab}|^2} \left[ E_{+}' E_{-}' - W_{aa}(E_{+}' + E_{-}') + |W_{ab}|^2 + W_{ab}^2 \right] \end{aligned}$$

en samment (6.27)

$$E_{+}' + E_{-}' = \frac{W_{aa} + W_{bb}}{2} \cdot 2$$

$$\text{eg } E'_+ E'_- = \frac{1}{4} \left\{ (W_{aa} + W_{bb})^2 - (W_{aa} - W_{bb})^2 - 4|W_{ab}|^2 \right\}$$

$$\rightarrow \langle \Phi_+^0 | \Phi_-^0 \rangle = \frac{\alpha_+^* \alpha_-}{|W_{ab}|^2} \left[ \frac{1}{4} (W_{aa} + W_{bb})^2 - \frac{1}{4} (W_{aa} - W_{bb})^2 - |W_{ab}|^2 \right. \\ \left. - W_{aa}^2 - W_{aa}W_{bb} + |W_{ab}|^2 + W_{aa}^2 \right]$$

$$= \frac{\alpha_+^* \alpha_-}{|W_{ab}|^2} \left[ \cancel{W_{aa}W_{bb}} - \cancel{W_{aa}^2} - \cancel{W_{aa}W_{bb}} + \cancel{W_{aa}^2} \right]$$

$$= 0$$

b) Syarat agar  $\langle \psi_+^0 | H' | \psi_-^0 \rangle = 0$

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$$\langle \psi_+^0 | H' | \psi_-^0 \rangle = \left\{ \alpha_+^* \langle \psi_a^0 | + \beta_+^* \langle \psi_b^0 | \right\} H' \left[ \alpha_- | \psi_a^0 \rangle + \beta_- | \psi_b^0 \rangle \right]$$

$$= \alpha_+^* \alpha_- \langle \psi_a^0 | H' | \psi_a^0 \rangle + \alpha_+^* \beta_- \langle \psi_a^0 | H' | \psi_b^0 \rangle$$

$$+ \beta_+^* \alpha_- \langle \psi_b^0 | H' | \psi_a^0 \rangle + \beta_+^* \beta_- \langle \psi_b^0 | H' | \psi_b^0 \rangle$$

$$= \alpha_+^* \alpha_- W_{aa} + \alpha_+^* \beta_- W_{ab} + \beta_+^* \alpha_- W_{ba} + \beta_+^* \beta_- W_{bb}$$

$$= \alpha_+^* \alpha_- \left[ W_{aa} + (E_-^i - W_{aa}) + \frac{(E_+^i - W_{aa}) W_{ba}}{W_{ab}^*} + \frac{(E_+^i - W_{aa}) (E_-^i - W_{aa}) W_{ab}}{W_{ab}^* W_{ab}} \right]$$

$$= \alpha_+^* \alpha_- \left[ W_{aa} + (E'_- - W_{aa}) + (E'_+ - W_{aa}) + \frac{(E'_+ - W_{aa})(E'_- - W_{aa})}{|W_{bb}|^2} W_{bb} \right] \quad (9)$$

frü (\*\*) setzt  $\Rightarrow (E'_+ - W_{aa})(E'_- - W_{aa}) = -|W_{bb}|^2$

$$\rightarrow \langle \psi_+^0 | H' | \psi_-^0 \rangle = \alpha_+^* \alpha_- \left[ W_{aa} + (E'_- - W_{aa}) + (E'_+ - W_{aa}) - W_{bb} \right]$$

$$= \alpha_+^* \alpha_- \left[ E'_- + E'_+ - W_{aa} - W_{bb} \right] = 0$$

↑  
Satzwert (6.27)

c) Symmetrie

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$$\langle \psi_{\pm}^0 | H' | \psi_{\pm}^0 \rangle = E'_{\pm}$$

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$$\langle \psi_{\pm}^0 | H' | \psi_{\pm}^0 \rangle = \left\{ \alpha_{\pm}^* \langle \psi_a^0 | + \beta_{\pm}^* \langle \psi_b^0 | \right\} H' \left\{ \alpha_{\pm} | \psi_a^0 \rangle + \beta_{\pm} | \psi_b^0 \rangle \right\}$$

$$= |\alpha_{\pm}|^2 W_{aa} + \alpha_{\pm}^* \beta_{\pm} W_{ab} + \beta_{\pm}^* \alpha_{\pm} W_{ba} + |\beta_{\pm}|^2 W_{bb}$$

$$= |\alpha_{\pm}|^2 \left\{ W_{aa} + \frac{(E'_{\pm} - W_{aa}) W_{ab}}{W_{ab}} \right\} + |\beta_{\pm}|^2 \left\{ W_{bb} + \frac{(E'_{\pm} - W_{bb}) W_{ba}}{W_{ba}} \right\}$$

$$= |\alpha_{\pm}|^2 \left\{ E'_{\pm} \right\} + |\beta_{\pm}|^2 \left\{ E'_{\pm} \right\} = \underline{E'_{\pm}}$$

$$= \{ |\alpha_{\pm}|^2 + |\beta_{\pm}|^2 \} E'_{\pm} = 1 \cdot E'_{\pm}$$