

①

Athugum  $(n, n-1, m) \rightarrow$  ástand vetisatans

①

Samkvæmt fjórirbætti og bók er almenna bylgjufallid

$$\psi_{n,m}(r) = \sqrt{\left(\frac{2}{na}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]}} e^{-\frac{r}{na}} \left(\frac{2r}{na}\right)^l L_{n-l-1}^{2l+1}\left(\frac{2r}{na}\right) Y_{lm}(\theta, \phi)$$

Ef  $l = n-1$ , (hæsta leyfilega  $l$ -gildi ástands með  $n$ )

þá er Laguerre  $L_0^{2n-1}\left(\frac{2r}{na}\right) = 1$

því  $2l+1 = 2n-1$  og  $n-l-1 = 0$

'Eg nota öðra skilgreiningu á Laguerre en bókun,

$$L_n^\alpha(x) = e^x \frac{x^{-\alpha}}{n!} \frac{d^n}{dx^n} (e^{-x} x^{n+\alpha})$$

$$\psi_{n,n-1,m}(r) = \sqrt{\left(\frac{2}{na}\right)^3 \frac{1}{2n[(2n-1)!]}} e^{-\frac{r}{na}} \left(\frac{2r}{na}\right)^{n-1} Y_{n-1,m}(\theta, \phi)$$

a) Athugum stöðumina

sjá

<http://en.wikipedia.org/Laguerre-polynomials>

öð. Lebedev. N.N., Dove

$$\begin{aligned}
& \int_0^{\infty} r^2 dr d\Omega |\psi_{n,n-1,m}(r)|^2 \\
&= \left(\frac{2}{na}\right)^3 \frac{1}{2n [(2n-1)!]} \int_0^{\infty} r^2 dr e^{-\frac{2r}{na}} \left(\frac{2r}{na}\right)^{2n-2} \\
&= \left(\frac{2}{na}\right)^3 \frac{n^3 a^3}{2n [(2n-1)!] 8} \int_0^{\infty} \left(\frac{2r}{na}\right)^2 d\left(\frac{2r}{na}\right) e^{-\frac{2r}{na}} \left(\frac{2r}{na}\right)^{2n-2} \\
&= \frac{1}{2n [(2n-1)!]} \int_0^{\infty} du u^{2n} e^{-u} = \frac{\Gamma(2n+1)}{2n [(2n-1)!]} = 1
\end{aligned}$$

b) Reikvid vertigildin fy —  $\langle r \rangle$  og  $\langle r^2 \rangle$

(3)

$$\langle r^p \rangle = \int_0^{\infty} r^{2+p} dr d\Omega |\psi_{n,n-1,m}(r)|^2$$

$$= \left(\frac{2}{na}\right)^3 \frac{1}{2^n (2n-1)!} \int_0^{\infty} r^{2+p} dr e^{-\frac{2r}{na}} \left(\frac{2r}{na}\right)^{2n-2}$$

$$= \left(\frac{2}{na}\right)^3 \frac{n^3 a^3}{2^n (2n-1)! \cdot 8} \frac{(na)^p}{2^p} \int_0^{\infty} d\left(\frac{2r}{na}\right) e^{-\frac{2r}{na}} \left(\frac{2r}{na}\right)^{2n} \left(\frac{2r}{na}\right)^p$$

$$= \left(\frac{na}{2}\right)^p \frac{1}{(2n)!} \int_0^{\infty} du e^{-u} u^{2n+p}$$

$$\langle r \rangle = \left(\frac{na}{2}\right) \frac{1}{(2n)!} \int_0^\infty du e^{-u} u^{2n+1} = \frac{na}{2} \frac{\Gamma(2n+2)}{(2n)!}$$

$$= \frac{na}{2} (2n+1) = na \left(n + \frac{1}{2}\right)$$

$$\langle r^2 \rangle = \left(\frac{na}{2}\right)^2 \frac{1}{(2n)!} \int_0^\infty du e^{-u} u^{2n+2} = \left(\frac{na}{2}\right)^2 \frac{\Gamma(2n+3)}{(2n)!}$$

$$= \left(\frac{na}{2}\right)^2 (2n+1)(2n+2) = (na)^2 \left(n + \frac{1}{2}\right)(n+1)$$

c)

$$\Delta_r = \sqrt{\langle r^2 \rangle - \langle r \rangle^2} = na \sqrt{\left\{ \left(n + \frac{1}{2}\right)(n+1) - \left(n + \frac{1}{2}\right)^2 \right\}}$$

(5)

$$\Delta_r = (na) \sqrt{\left(n + \frac{1}{2}\right)^2} \cdot \frac{1}{\sqrt{2}}$$

$$\rightarrow \Delta_r / \langle r \rangle = \frac{1}{\sqrt{2}} \frac{1}{\sqrt{n + \frac{1}{2}}} = \frac{1}{\sqrt{2n+1}}$$

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$\langle r \rangle \sim n^2 a \rightarrow$  fjarlægð rafendur vex sem  $n^2$

en  $\frac{\Delta_r}{\langle r \rangle} = \frac{1}{\sqrt{2n+1}}$  þýðir að bylgjufallid

þengist með vaxandi  $n$

$\rightarrow$  stefni á klassíska breyt

d) Teikna fyrir  $n = 1, 5, 17$

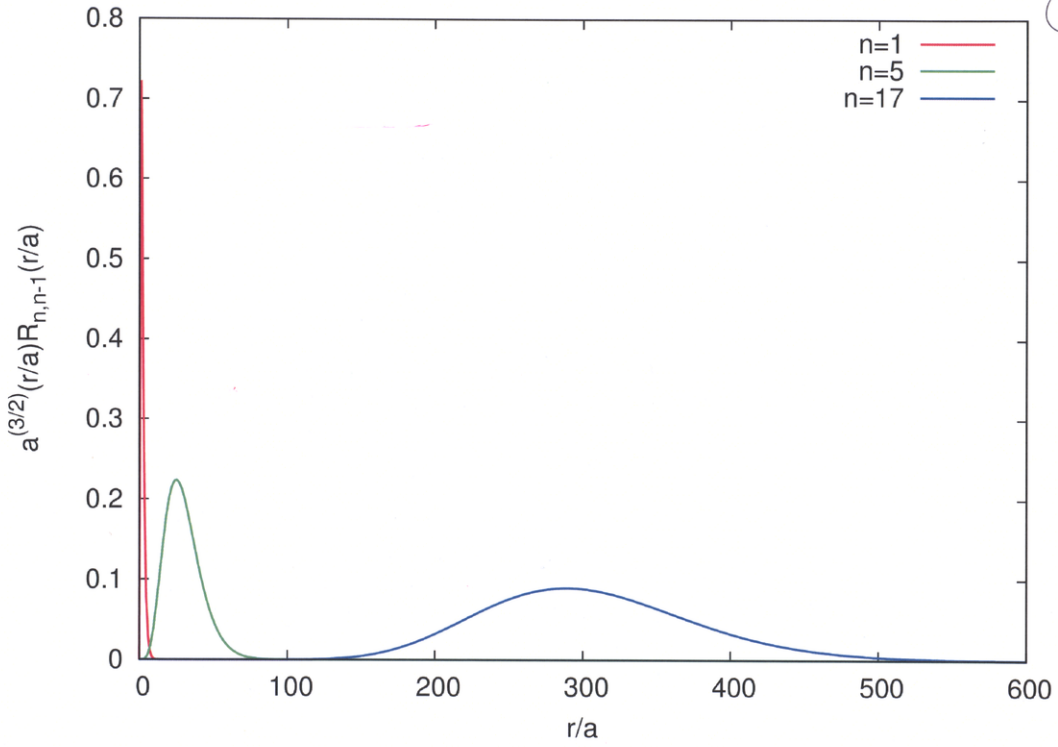
(6)

$$R_{n,n-1}(r) = \sqrt{\left(\frac{2}{na}\right)^3 \frac{1}{2n(2n-1)!}} e^{-\frac{r}{na}} \left(\frac{2r}{na}\right)^{n-1}$$

$$R_{n,n-1}\left(\frac{r}{a}\right) = \sqrt{\left(\frac{2}{na}\right)^3 \frac{1}{2n(2n-1)!}} e^{-\frac{1}{n}\left(\frac{r}{a}\right)} \left(\frac{r}{a}\right)^{n-1} \cdot \left(\frac{2}{n}\right)^{n-1}$$

$$a^{3/2} \left(\frac{r}{a}\right) R_{n,n-1}\left(\frac{r}{a}\right) = \left(\frac{2}{n}\right)^{n-1} \sqrt{\left(\frac{2}{n}\right)^3 \frac{1}{2n(2n-1)!}} e^{-\frac{1}{n}\left(\frac{r}{a}\right)} \left(\frac{r}{a}\right)^n$$

$$a^{3/2} \left(\frac{r}{a}\right) R_{n,n-1}\left(\frac{r}{a}\right) = F(n) e^{-\frac{1}{n}\left(\frac{r}{a}\right)} \left(\frac{r}{a}\right)^n$$



② Finna fylkja-útselningu  $S_y$  og  $S_z$  fyrir  
eind með spuna  $\frac{3}{2}$  í grunni eiginástanda  $S_z$  ⑧

$S_z$  hefur eiginástandin

$$\left| \frac{3}{2}, +\frac{3}{2} \right\rangle \quad \left| \frac{3}{2}, +\frac{1}{2} \right\rangle \quad \left| \frac{3}{2}, -\frac{1}{2} \right\rangle \quad \left| \frac{3}{2}, -\frac{3}{2} \right\rangle$$

Notum (4.136):  $S_{\pm} |s, m\rangle = \hbar \sqrt{s(s+1) - m(m\pm 1)} |s, m\pm 1\rangle$

$$\text{og } S_y = \frac{1}{2i} (S_+ - S_-)$$

$$S_+ \left| \frac{3}{2}, +\frac{3}{2} \right\rangle = 0$$

$$S_+ \left| \frac{3}{2}, +\frac{1}{2} \right\rangle = \hbar \sqrt{\frac{3}{2} \cdot \frac{5}{2} - \frac{1}{2} \cdot \frac{3}{2}} \left| \frac{3}{2}, +\frac{3}{2} \right\rangle = \sqrt{3} \hbar \left| \frac{3}{2}, +\frac{3}{2} \right\rangle$$



$$S_+ \left| \frac{3}{2}, -\frac{1}{2} \right\rangle = 2\hbar \left| \frac{3}{2}, +\frac{1}{2} \right\rangle \quad \left| S_- \left| \frac{3}{2}, +\frac{3}{2} \right\rangle = \sqrt{3}\hbar \left| \frac{3}{2}, +\frac{1}{2} \right\rangle \right. \quad (9)$$

$$S_+ \left| \frac{3}{2}, -\frac{3}{2} \right\rangle = \sqrt{3}\hbar \left| \frac{3}{2}, -\frac{1}{2} \right\rangle \quad S_- \left| \frac{3}{2}, +\frac{1}{2} \right\rangle = 2\hbar \left| \frac{3}{2}, -\frac{1}{2} \right\rangle$$

$$S_- \left| \frac{3}{2}, -\frac{1}{2} \right\rangle = \sqrt{3}\hbar \left| \frac{3}{2}, -\frac{3}{2} \right\rangle \quad S_- \left| \frac{3}{2}, -\frac{3}{2} \right\rangle = 0$$

$$\rightarrow S_+ = \hbar \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & \sqrt{3} \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$S_- = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ \sqrt{3} & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}$$

$$S_y = \frac{1}{2i} (S_+ - S_-)$$

$$= \frac{\hbar}{2i} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ -\sqrt{3} & 0 & 2 & 0 \\ 0 & -2 & 0 & \sqrt{3} \\ 0 & 0 & -\sqrt{3} & 0 \end{pmatrix}$$

Eigengildi  $S_y$

10

Reynt ~~með~~ octave getur

$$+ \frac{h}{2} \cdot 3$$

eins og ~~er~~ matthé búa

$$+ \frac{h}{2} \cdot 1$$

$$- \frac{h}{2} \cdot 1$$

$$- \frac{h}{2} \cdot 3$$